<u>Electric Potential Function</u> <u>for Charge Densities</u>

Recall the total static electric field produced by 2 **different** charges (or charge densities) is just the **vector sum** of the fields produced by each:

$$\boldsymbol{\mathsf{E}}\left(\overline{\boldsymbol{\mathsf{r}}}\right) = \boldsymbol{\mathsf{E}}_{\!\!1}\left(\overline{\boldsymbol{\mathsf{r}}}\right) + \boldsymbol{\mathsf{E}}_{\!\!2}\left(\overline{\boldsymbol{\mathsf{r}}}\right)$$

Since the fields are conservative, we can write this as:

$$\mathbf{E}(\overline{\mathbf{r}}) = \mathbf{E}_{1}(\overline{\mathbf{r}}) + \mathbf{E}_{2}(\overline{\mathbf{r}})$$
$$-\nabla \mathbf{V}(\overline{\mathbf{r}}) = -\nabla \mathbf{V}_{1}(\overline{\mathbf{r}}) - \nabla \mathbf{V}_{2}(\overline{\mathbf{r}})$$
$$-\nabla \mathbf{V}(\overline{\mathbf{r}}) = -\nabla \left(\mathbf{V}_{1}(\overline{\mathbf{r}}) + \mathbf{V}_{2}(\overline{\mathbf{r}})\right)$$

Therefore, we find,

 $V(\overline{r}) = V_1(\overline{r}) + V_2(\overline{r})$

In other words, **superposition** also holds for the electric potential function! The total electric potential field produced by a collection of charges is simply the **sum** of the electric potential produced by **each**.

Consider now some distribution of charge, $\rho_v(\bar{r})$. The amount of charge dQ, contained within small volume dv, located at position \bar{r}' , is: $dQ = \rho_v(\bar{r}') dv'$ The **electric potential function** produced by this charge is therefore:

$$dV(\bar{\mathbf{r}}) = \frac{dQ}{4\pi\varepsilon_0 |\bar{\mathbf{r}}-\bar{\mathbf{r}}'|}$$
$$= \frac{\rho_v(\bar{\mathbf{r}}') dv'}{4\pi\varepsilon_0 |\bar{\mathbf{r}}-\bar{\mathbf{r}}'|}$$

Therefore, **integrating** across all the charge in some **volume** V, we get:

$$\mathcal{V}\left(\overline{\mathbf{r}}\right) = \iiint_{\mathcal{V}} \frac{\rho_{\nu}\left(\overline{\mathbf{r}}'\right)}{4\pi\varepsilon_{0}\left|\overline{\mathbf{r}}-\overline{\mathbf{r}}'\right|} d\nu'$$

Likewise, for **surface** or **line** charge density:

$$V(\overline{\mathbf{r}}) = \iint_{S} \frac{\rho_{s}(\overline{\mathbf{r}}')}{4\pi\varepsilon_{0} |\overline{\mathbf{r}} - \overline{\mathbf{r}}'|} ds'$$

$$V(\overline{\mathbf{r}}) = \int_{\mathcal{C}} \frac{\rho_{\ell}(\overline{\mathbf{r}}')}{4\pi\varepsilon_{0} |\overline{\mathbf{r}} - \overline{\mathbf{r}}'|} d\ell'$$

Note that these integrations are **scalar** integrations—typically they are **easier** to evaluate than the integrations resulting from **Coulomb's Law**.

Once we find the electric potential function $V(\bar{r})$, we can **then** determine the total **electric field** by taking the gradient:

$$\mathsf{E}(\bar{r}) = -\nabla \mathsf{V}(\bar{r})$$

Thus, we now have three (!) potential methods for determining the electric field produced by some charge distribution $\rho_{\nu}(\bar{r})$:

- 1. Determine $E(\overline{r})$ from Coulomb's Law.
- 2. If $\rho_v(\bar{r})$ is symmetric, determine $\mathbf{E}(\bar{r})$ from Gauss's Law.
- 3. Determine the electric potential function $V(\bar{r})$, and then determine the electric field as $E(\bar{r}) = -\nabla V(\bar{r})$.

Q: Yikes! Which of the three should we use??

A: To a certain extent, it does **not matter**! All three will provide the **same** result (although $\rho_{\nu}(\bar{r})$ **must** be symmetric to use method 2!).

However, **if** the charge density is symmetric, we will find that using Gauss's Law (method 2) will **typically** result in much less work!

Otherwise (i.e., for **non**-symmetric $\rho_v(\bar{r})$), we find that **sometimes** method 1 is easiest, but in **other** cases method 3 is a bit less stressful (i.e., **you** decide!).