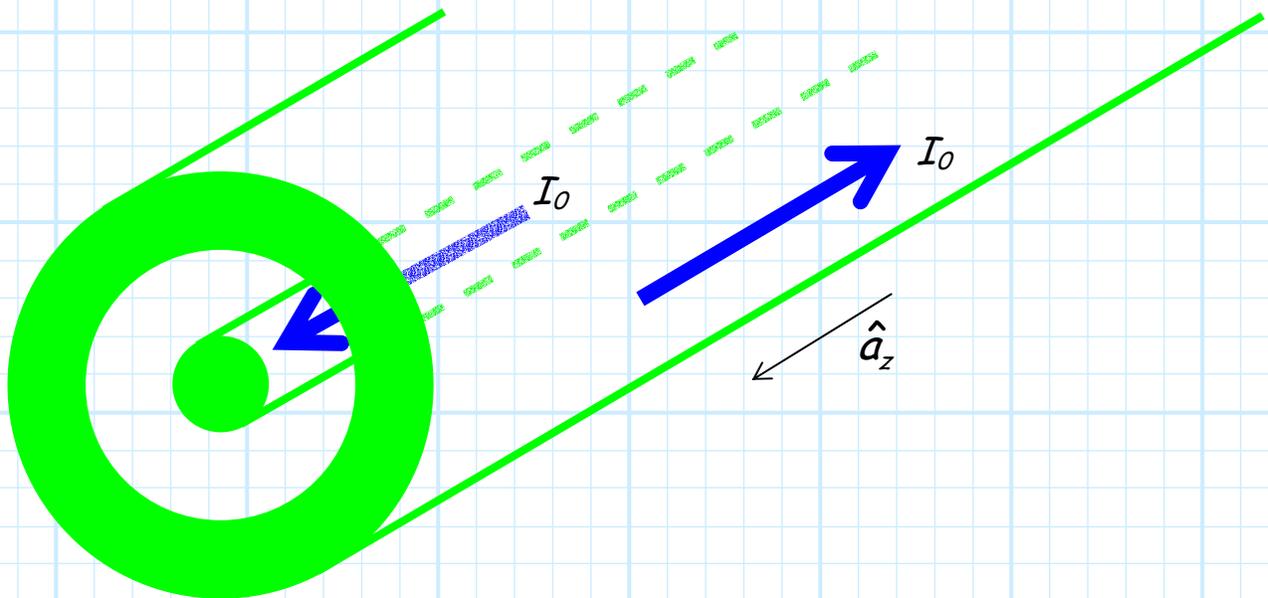
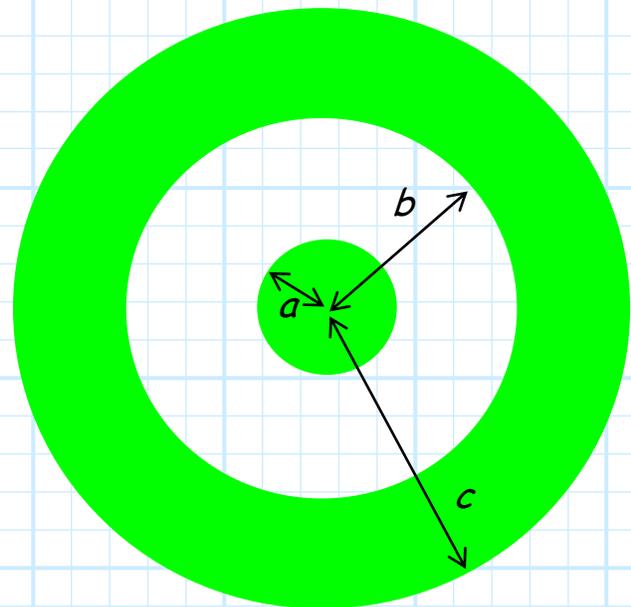


# Example: The B-Field of Coaxial Transmission Line

Consider now a coaxial cable, with inner radius  $a$ :



The **outer** surface of the **inner** conductor has radius  $a$ , the **inner** surface of the **outer** conductor has radius  $b$ , and the **outer** radius of the **outer** conductor has radius  $c$ .



Typically, the **current** flowing on the **inner** conductor is **equal but opposite** that flowing in the **outer** conductor. Thus, if current  $I_0$  is flowing in the **inner** conductor in the direction  $\hat{a}_z$ , then current  $I_0$  will be flowing in the **outer** conductor in the opposite (i.e.,  $-\hat{a}_z$ ) direction.

**Q:** *Hey! If there is current, a magnetic flux density must be created. What is the vector field  $\mathbf{B}(\bar{r})$ ?*

**A:** We've already determined this (sort of)!

Recall we found the magnetic flux density produced by a hollow cylinder—we can use this to determine the magnetic flux density in a coaxial transmission line.

**→** A coaxial cable can be viewed as **two** hollow cylinders!



**Q:** *I find it necessary to point out that you are indeed **wrong**—the **inner** conductor is **not hollow**!*

**A:** Mathematically, we can view the inner conductor as a hollow cylinder with an outer radius  $a$  and an **inner** radius of **zero**!

Thus, we can use the results of the previous handout to conclude that the magnetic flux density produced by the current flowing in the inner conductor is:

$$\mathbf{B}_{inner}(\bar{r}) = \begin{cases} \frac{I_0 \mu_0}{2\pi\rho} \left( \frac{\rho^2 - 0^2}{a^2 - 0^2} \right) \hat{a}_\phi = \frac{I_0 \mu_0}{2\pi a^2} \rho \hat{a}_\phi & \rho < a \\ \frac{I_0 \mu_0}{2\pi\rho} \hat{a}_\phi & \rho > a \end{cases} \quad \left[ \frac{\text{Webers}}{m^2} \right]$$

Likewise, we can use the same result to determine the magnetic flux density of the current flowing in the outer conductor:

$$\mathbf{B}_{outer}(\bar{r}) = \begin{cases} 0 & \rho < b \\ \frac{-I_0 \mu_0}{2\pi\rho} \left( \frac{\rho^2 - b^2}{c^2 - b^2} \right) \hat{a}_\phi & b < \rho < c \\ \frac{-I_0 \mu_0}{2\pi\rho} \hat{a}_\phi & \rho > c \end{cases} \quad \left[ \frac{\text{Webers}}{m^2} \right]$$

Note the **minus sign** is due to direction of the current ( $-\hat{a}_z$ ) in the outer conductor.

We can now apply **superposition** to determine the total magnetic flux density in a coaxial transmission line! Specifically:

$$\text{if } \mathbf{J}(\bar{\mathbf{r}}) = \mathbf{J}_{inner}(\bar{\mathbf{r}}) + \mathbf{J}_{outer}(\bar{\mathbf{r}})$$

$$\text{then } \mathbf{B}(\bar{\mathbf{r}}) = \mathbf{B}_{inner}(\bar{\mathbf{r}}) + \mathbf{B}_{outer}(\bar{\mathbf{r}})$$

Note due to the **piecewise** nature of these solutions, we must evaluate this sum for **4** distinct regions:

- 1)  $\rho < a$  (in the inner conductor)
- 2)  $a < \rho < b$  (in the region between the conductors)
- 3)  $b < \rho < c$  (in the outer conductor)
- 4)  $\rho > c$  (outside the coaxial cable)

$$\underline{\rho < a}$$

$$\begin{aligned} \mathbf{B}(\bar{\mathbf{r}}) &= \mathbf{B}_{inner}(\bar{\mathbf{r}}) + \mathbf{B}_{outer}(\bar{\mathbf{r}}) \\ &= \frac{I_0 \mu_0}{2\pi a^2} \rho \hat{\mathbf{a}}_\phi + 0 \\ &= \frac{I_0 \mu_0}{2\pi a^2} \rho \hat{\mathbf{a}}_\phi \end{aligned}$$

$$\underline{a < \rho < b}$$

$$\begin{aligned} \mathbf{B}(\bar{\mathbf{r}}) &= \mathbf{B}_{inner}(\bar{\mathbf{r}}) + \mathbf{B}_{outer}(\bar{\mathbf{r}}) \\ &= \frac{I_0 \mu_0}{2\pi \rho} \hat{\mathbf{a}}_\phi + 0 \\ &= \frac{I_0 \mu_0}{2\pi \rho} \hat{\mathbf{a}}_\phi \end{aligned}$$

$$\underline{b < \rho < c}$$

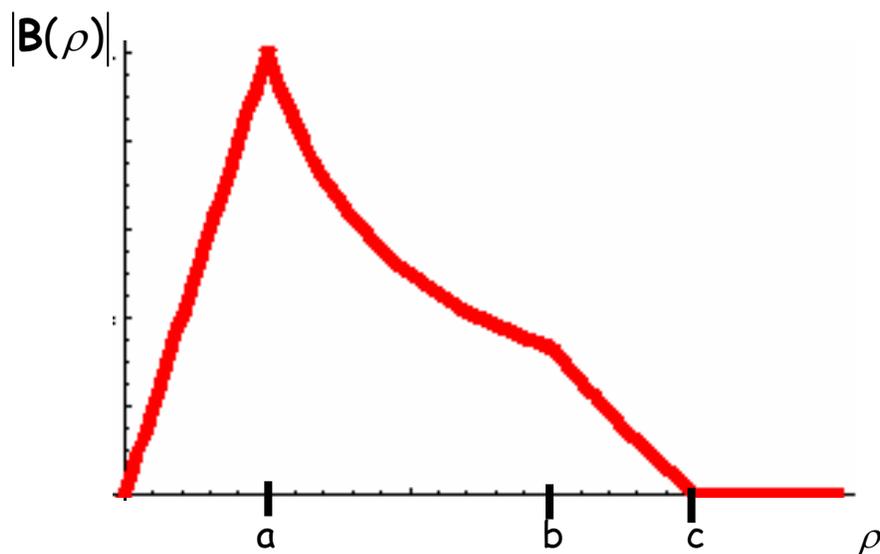
$$\begin{aligned} \mathbf{B}(\bar{\mathbf{r}}) &= \mathbf{B}_{inner}(\bar{\mathbf{r}}) + \mathbf{B}_{outer}(\bar{\mathbf{r}}) \\ &= \frac{I_0 \mu_0}{2\pi \rho} \hat{\mathbf{a}}_\phi + \frac{-I_0 \mu_0}{2\pi \rho} \left( \frac{\rho^2 - b^2}{c^2 - b^2} \right) \hat{\mathbf{a}}_\phi \\ &= \frac{I_0 \mu_0}{2\pi \rho} \left( \frac{c^2 - \rho^2}{c^2 - b^2} \right) \hat{\mathbf{a}}_\phi \end{aligned}$$

$$\underline{\rho > c}$$

$$\begin{aligned} \mathbf{B}(\bar{\mathbf{r}}) &= \mathbf{B}_{inner}(\bar{\mathbf{r}}) + \mathbf{B}_{outer}(\bar{\mathbf{r}}) \\ &= \frac{I_0 \mu_0}{2\pi \rho} \hat{\mathbf{a}}_\phi + \frac{-I_0 \mu_0}{2\pi \rho} \hat{\mathbf{a}}_\phi \\ &= 0 \end{aligned}$$

Summarizing, we find the total magnetic flux density to be:

$$\mathbf{B}(\bar{r}) = \begin{cases} \frac{I_0 \mu_0}{2\pi a^2} \rho \hat{a}_\phi & \rho < a \\ \frac{I_0 \mu_0}{2\pi \rho} \hat{a}_\phi & a < \rho < b \\ \frac{I_0 \mu_0}{2\pi \rho} \left( \frac{c^2 - \rho^2}{c^2 - b^2} \right) \hat{a}_\phi & b < \rho < c \\ 0 & \rho > c \end{cases} \quad \left[ \frac{\text{Webers}}{m^2} \right]$$



The magnetic flux density **and** the electric field **outside** of a coaxial transmission line are **zero!**