Magnetization Currents

Recall that the magnetic vector potential $\mathbf{A}(\bar{r})$ created by volume current distribution $\mathbf{J}(\bar{r})$ is:

$$\mathbf{A}(\overline{\mathbf{r}}) = \frac{\mu_0}{4\pi} \iiint_{\mathbf{v}} \frac{\mathbf{J}(\overline{\mathbf{r}}')}{|\overline{\mathbf{r}} - \overline{\mathbf{r}}'|} d\mathbf{v}'$$

while the magnetic vector potential created by a surface current $\mathbf{J}_{s}(\overline{r})$:

$$\mathbf{A}(\overline{\mathbf{r}}) = \frac{\mu_0}{4\pi} \iint_{\mathcal{S}} \frac{\mathbf{J}_s(\overline{\mathbf{r}}')}{|\overline{\mathbf{r}} - \overline{\mathbf{r}}'|} ds'$$

Therefore, if **both** volume and surface current densities are present we find that the **total** magnetic vector potential is:

$$\boldsymbol{A}(\overline{\mathbf{r}}) = \frac{\mu_0}{4\pi} \iiint_{\nu} \frac{\mathbf{J}(\overline{\mathbf{r}}')}{|\overline{\mathbf{r}} - \overline{\mathbf{r}}'|} d\nu' + \frac{\mu_0}{4\pi} \iint_{\mathcal{S}} \frac{\mathbf{J}_s(\overline{\mathbf{r}}')}{|\overline{\mathbf{r}} - \overline{\mathbf{r}}'|} ds'$$

Compare these expressions to the magnetic vector potential field produced by material with **Magnetization Vector** $M(\bar{r})$:

$$\mathbf{A}(\overline{\mathbf{r}}) = \iiint_{\nu} \frac{\mu_0 \mathbf{M}(\overline{\mathbf{r}}') \times (\overline{\mathbf{r}} - \overline{\mathbf{r}}')}{4\pi |\overline{\mathbf{r}} - \overline{\mathbf{r}}'|^3} d\nu$$

We can write also write this expression as (trust me!):

$$\mathbf{A}(\overline{\mathbf{r}}) = \frac{\mu_0}{4\pi} \iiint_{\mathbf{v}} \frac{\nabla' \mathbf{x} \mathbf{M}(\overline{\mathbf{r}}')}{|\overline{\mathbf{r}} - \overline{\mathbf{r}}'|} d\mathbf{v}' + \frac{\mu_0}{4\pi} \bigoplus_{\mathbf{s}} \frac{\mathbf{M}(\overline{\mathbf{r}}') \mathbf{x} \hat{\mathbf{a}}_{\mathbf{n}}}{|\overline{\mathbf{r}} - \overline{\mathbf{r}}'|} d\mathbf{s}'$$

where surface S is the **closed surface** that surrounds material volume V, and unit vector \hat{a}_n is **normal** to this surface.

We find that this is identical to the expression:

$$\mathbf{A}(\overline{\mathbf{r}}) = \frac{\mu_0}{4\pi} \iiint_{\mathbf{V}} \frac{\mathbf{J}(\overline{\mathbf{r}}')}{|\overline{\mathbf{r}} - \overline{\mathbf{r}}'|} \, d\mathbf{v}' + \frac{\mu_0}{4\pi} \iint_{\mathbf{S}} \frac{\mathbf{J}_{\mathbf{s}}(\overline{\mathbf{r}}')}{|\overline{\mathbf{r}} - \overline{\mathbf{r}}'|} \, d\mathbf{s}$$

if $\mathbf{J}(\bar{r}) = \nabla \mathbf{x} \mathbf{M}(\bar{r})$ and $\mathbf{J}_{s}(\bar{r}) = \mathbf{M}(\bar{r}) \mathbf{x} \hat{a}_{n}$.

Therefore, we find that the magnetization of some material, as described by magnetization vector $\mathbf{M}(\overline{r})$, creates **effective** currents $\mathbf{J}_m(\overline{r})$ and $\mathbf{J}_{sm}(\overline{r_s})$ (where $\overline{r_s}$ indicates points on the material surface). We call these effective currents magnetization currents:

$$\mathbf{J}_{m}(\bar{\boldsymbol{r}}) = \nabla \mathbf{x} \, \mathbf{M}(\bar{\boldsymbol{r}}) \qquad \left[\frac{\boldsymbol{A}}{m^{2}}\right]$$

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 $\mathbf{J}_{sm}\left(\overline{r_{s}}\right) = \mathbf{M}\left(\overline{r_{s}}\right) \times \mathbf{\hat{a}}_{n}$

Again, note the **analogy** of these **magnetization** currents with **polarization** charges $\rho_{vp}(\bar{r})$ and $\rho_{sp}(\bar{r})$.