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6-2 Capacitance

Reading Assignment: pp. 179-185

To **create** a potential difference *V* between two perfect conductors, we must **deposit** charge Q on the conductor surface.

Or, if we **enforce** a potential difference *V* (with a voltage source) between the two perfect conductors, some charge Q will be **stored** on the conductor surface.

Q:

A:

HO: Capacitance

HO: The Parallel Plate Capacitor

HO: Capacitance of a Coaxial Transmission Line

<u>Capacitance</u>

Consider two conductors, with a potential difference of V volts.

 V_0

E(r

* Since there is a potential difference between the conductors, there must be an **electric potential field** $V(\bar{r})$, and therefore an **electric field** $E(\bar{r})$ in the region between the conductors.

* Likewise, if there is an electric field, then we can specify an **electric flux density** $D(\overline{r})$, which we can use to determine the **surface charge density** $\rho_s(\overline{r})$ on each of the conductors.

* We find that if the total net charge on **one** conductor is Q then the charge on the **other** will be equal to -Q.

In other words, the total net charge on each conductor will be **equal** but **opposite**!

 $-\rho_{\epsilon}$

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Note that this does **not** mean that the surface charge densities on each conductor are equal (i.e., $\rho_{s_+}(\overline{r}) \neq \rho_{s_-}(\overline{r})$). Rather, it means that:

 $\bigoplus_{S_{+}} \rho_{s+}(\overline{r}) ds = - \bigoplus_{S_{-}} \rho_{s-}(\overline{r}) ds = Q$

where surface S_+ is the surface surrounding the conductor with the positive charge (and the higher electric potential), while the surface S_- surrounds the conductor with the negative charge.

Q: How much free **charge** Q is there on each conductor, and how does this charge relate to the **voltage** V_0 ?

A: We can determine this from the mutual capacitance C of these conductors!

The mutual **capacitance** between two conductors is **defined** as:

$$C = \frac{Q}{V} \qquad \left[\frac{Coulombs}{Volt} \doteq Farad \right]$$

where Q is the total charge on each conductor, and *V* is the **potential difference** between each conductor (for our example, $V = V_0$).

Recall that the total charge on a conductor can be determined by **integrating** the surface charge density $\rho_s(\overline{r})$ across the **entire surface** S of a conductor:

$$Q = \bigoplus_{c} \rho_{s+}(\overline{r}) ds = -\bigoplus_{c} \rho_{s-}(\overline{r}) ds$$

But recall also that the surface charge density on the surface of a conductor can be determined from the **electric flux** density $D(\bar{r})$:

$$\rho_{s}\left(\overline{\mathbf{r}}\right) = \mathbf{D}\left(\overline{\mathbf{r}}\right) \cdot \hat{a}_{n}$$

where \hat{a}_n is a unit vector **normal** to the conductor.

Combining the two equations above, we get:

$$Q = \bigoplus_{S_{+}} \mathbf{D}(\overline{\mathbf{r}}) \cdot \hat{a}_{n} ds = -\bigoplus_{S_{-}} \mathbf{D}(\overline{\mathbf{r}}) \cdot \hat{a}_{n} ds$$
$$= \bigoplus_{S_{+}} \mathbf{D}(\overline{\mathbf{r}}) \cdot \overline{ds} = -\bigoplus_{S_{-}} \mathbf{D}(\overline{\mathbf{r}}) \cdot \overline{ds}$$

where we remember that $\overline{ds} = \hat{a}_n ds$.

Hey! This is **no surprise**! We **already** knew that:

$$Q = \bigoplus D(\overline{r}) \cdot \overline{ds}$$

This expression is also know as

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Note since $D(\overline{r}) = \varepsilon E(\overline{r})$ we can also say:

$$Q = \bigoplus_{r} \varepsilon \mathbf{E}(\overline{r}) \cdot \overline{ds}$$

The **potential difference** *V* between two conductors can likewise be determined as:

$$V = \int \mathbf{E}(\overline{\mathbf{r}}) \cdot \overline{d\ell}$$

where C is any contour that leads from one conductor to the other.

Q: Why any contour?

A:

We can therefore determine the **capacitance** between two conductors as:

$$C = \frac{Q}{V} = \frac{\bigoplus_{S_{+}} \varepsilon \mathbf{E}(\overline{\mathbf{r}}) \cdot \overline{ds}}{\int_{C} \mathbf{E}(\overline{\mathbf{r}}) \cdot \overline{d\ell}} \quad [Farad]$$

Where the contour C must start at some point on surface S_+ and end at some point on surface S_- .

Note this expression can be written as:

$$Q = C V$$

In other words, the charge **stored** by two conductors is equal to the product of their mutual capacitance and the potential difference between them.

Therefore, the greater capacitance, the greater the amount of charge that is stored.

By the way, try taking the **time derivative** of the above equation:

$$\frac{dQ}{dt} = C \frac{dV}{dt}$$
$$I = C \frac{dV}{dt}$$

Look familiar ?

By the way, the current *I* in this equation is **displacement** current.

<u>The Parallel</u> <u>Plate Capacitor</u>

Consider the geometry of a **parallel plate capacitor**:



to approximate the fields and charge densities of this **finite** structure, where the **area** of each plate is 5.

For example, we determined that the **surface charge density** on the upper plate is:

$$\rho_{s+}\left(\overline{\mathbf{r}}\right) = \frac{\varepsilon V_0}{d}$$

The total charge on the upper plate is therefore:



The **capacitance** of this structure is therefore:

$$C = \frac{Q}{V} = \left(\frac{\varepsilon V_0 S}{d}\right) \left(\frac{1}{V_0}\right) = \frac{\varepsilon S}{d}$$
 [Farads]

Note therefore, that we can **increase** the capacitance of a parallel plate capacitor by:



Q: What is the **capacitance** between these two conducting structures?

A: The potential difference between them is V_0 . The **total charge** on one conducting structure is simply the sum of the charges on each plate:



<u>Capacitance of a Coaxial</u> <u>Transmission Line</u>

Recall the geometry of a coaxial transmission line:



The total charge Q on the inner conductor of a coax of length ℓ is determined by integrating the surface charge density across the conductor surface:

$$Q = \bigoplus_{S_{+}}^{\ell} \rho_{sa}(\bar{r}) ds$$
$$= \int_{0}^{\ell} \int_{0}^{2\pi} \frac{\varepsilon V_{0}}{\ln[b/a]} \frac{1}{a} \rho d\phi dz$$
$$= \frac{\varepsilon V_{0}}{\ln[b/a]} \frac{1}{a} \int_{0}^{\ell} \int_{0}^{2\pi} a d\phi dz$$
$$= \frac{\varepsilon V_{0}}{\ln[b/a]} \frac{1}{a} \ell(2\pi a)$$
$$= V_{0} \frac{2\pi \varepsilon}{\ln[b/a]} \ell$$

Note since $\rho_{sa}(\overline{r}) = \mathbf{D}(\overline{r}) \cdot \hat{a}_n$, we would have arrived at the same result by using:

$$Q = \oint \varepsilon \mathbf{E}(\overline{\mathbf{r}}) \cdot \overline{ds}$$

We can now determine the capacitance of this coaxial line!

Since C = Q/V, and since the **potential difference** between the conductors is $V = V_0$, we find:

$$C = \frac{Q}{V} = \left(V_0 \frac{2\pi \varepsilon}{\ln[b/a]}\ell\right) \left(\frac{1}{V_0}\right)$$
$$= \frac{2\pi \varepsilon}{\ln[b/a]}\ell$$

This value represents the capacitance of a coaxial line of length ℓ . A more useful expression is the capacitance of a coaxial line **per unit length** (e.g. farads/meter). We find this simply by **dividing** by length ℓ :

 $\frac{\mathcal{C}}{\ell} = \frac{2\pi \varepsilon}{\ln \lfloor b/a \rfloor} \qquad \left[\frac{\text{farads}}{\text{meter}} \right]$

Note the longer the transmission line, the greater the capacitance!

This can cause **great difficulty** if the voltage across the transmission line conductors is **time varying** (as it almost certainly will be!).

For **long** transmission lines, engineers cannot consider a transmission line simply as a "**wire**" conductor that connects circuit elements together. Instead, capacitance (and inductance) make the transmission line **itself** a **circuit element**!

In this case, engineers must use **transmission line theory** to analyze circuits !