

# The Stored Energy of Charge Distributions

Consider the case where a 1 Coulomb **point charge** is located at the **origin**. A **second charge  $Q$**  is **moved** to a distance  $r$  from the origin.

**Q:** *How much **energy** is stored in this simple charge distribution?*

**A:** Precisely the amount of **work** required to construct it!

Recall the amount of **work** required to move a charge  $Q$  through an electric field is:

$$W = -Q \int_c \mathbf{E}(\bar{r}) \cdot \overline{d\ell}$$

The work required to **move** the a charge from **infinity** to a distance  $r$  from the origin is therefore:

$$W = -Q \int_{\infty}^r \mathbf{E}(\bar{r}) \cdot \hat{a}_r dr$$

The 1 Coulomb charge at the origin of course produces the **electric field**:

$$\mathbf{E}(\bar{r}) = \frac{\hat{a}_r}{4\pi\epsilon r^2}$$

And produces the electric **potential** field:

$$V(\bar{r}) = \frac{1}{4\pi\epsilon r}$$

The work required to move charge  $Q$  to a distance  $r$  in **these** fields is:

$$\begin{aligned} W &= -Q \int_{\infty}^r \mathbf{E}(\bar{r}) \cdot \hat{a}_r dr \\ &= -Q \int_{\infty}^r \frac{1}{4\pi\epsilon r^2} \hat{a}_r \cdot \hat{a}_r dr \\ &= -\frac{Q}{4\pi\epsilon} \int_{\infty}^r \frac{1}{r^2} dr \\ &= \frac{Q}{4\pi\epsilon} \frac{1}{r} \end{aligned}$$

If we examine this result, we see that it is simply the **product** of the **charge**  $Q$  and the electric **potential** field  $V(\bar{r})$ .

$$\begin{aligned} W &= \frac{Q}{4\pi\epsilon} \frac{1}{r} \\ &= Q \frac{1}{4\pi\epsilon r} \\ &= Q V(\bar{r}) \end{aligned}$$

This seems to make sense! The units of electric potential are **Joules/Coulomb**, and the units of charge are of course **Coulombs**. The product of these two is therefore **energy**.

For a more general case, we find the **work** required to construct a **charge distribution**  $\rho_v(\bar{r})$  is:

$$W_e = \frac{1}{2} \iiint_V \rho_v(\bar{r}) V(\bar{r}) dV$$

This equation, therefore, is **also** equal to the (potential) **energy stored** by this charge distribution!

$W_e =$  potential **energy** stored by a charge distribution

Recall that charge density is related to electric flux density via the point form of Gauss's Law:

$$\nabla \cdot \mathbf{D}(\bar{r}) = \rho_v(\bar{r})$$

Likewise, the electric field is related to the electric potential as:

$$\mathbf{E}(\bar{r}) = -\nabla V(\bar{r})$$

As shown on page 198, we can use these expressions to rewrite the stored energy in terms of the electric field and the electric flux density:

$$W_e = \frac{1}{2} \iiint_V \mathbf{D}(\bar{\mathbf{r}}) \cdot \mathbf{E}(\bar{\mathbf{r}}) \, dv$$

What these expressions mean is that it takes energy to assemble a charge distribution  $\rho_v(\bar{\mathbf{r}})$ , or equivalently, an electric field  $\mathbf{E}(\bar{\mathbf{r}})$ . This energy is stored until it is released—the charge density returns to zero.

**Q:** *Is this energy stored in the fields  $\mathbf{E}(\bar{\mathbf{r}})$  and  $\mathbf{D}(\bar{\mathbf{r}})$ , or by the charge  $\rho_v(\bar{\mathbf{r}})$ ??*

**A:** One equation for  $W_e$  would suggest that the energy is stored by the **fields**, while the other by the **charge**.

It turns out, **either** interpretation is correct! The fields  $\mathbf{E}(\bar{\mathbf{r}})$  and  $\mathbf{D}(\bar{\mathbf{r}})$  **cannot** exist without a charge density  $\rho_v(\bar{\mathbf{r}})$ , and knowledge of the fields allow us to determine **completely** the charge density.

In other words, charges and the fields they create are “inseparable pairs”, since both must be present, we can attribute the stored energy to **either** quantity.