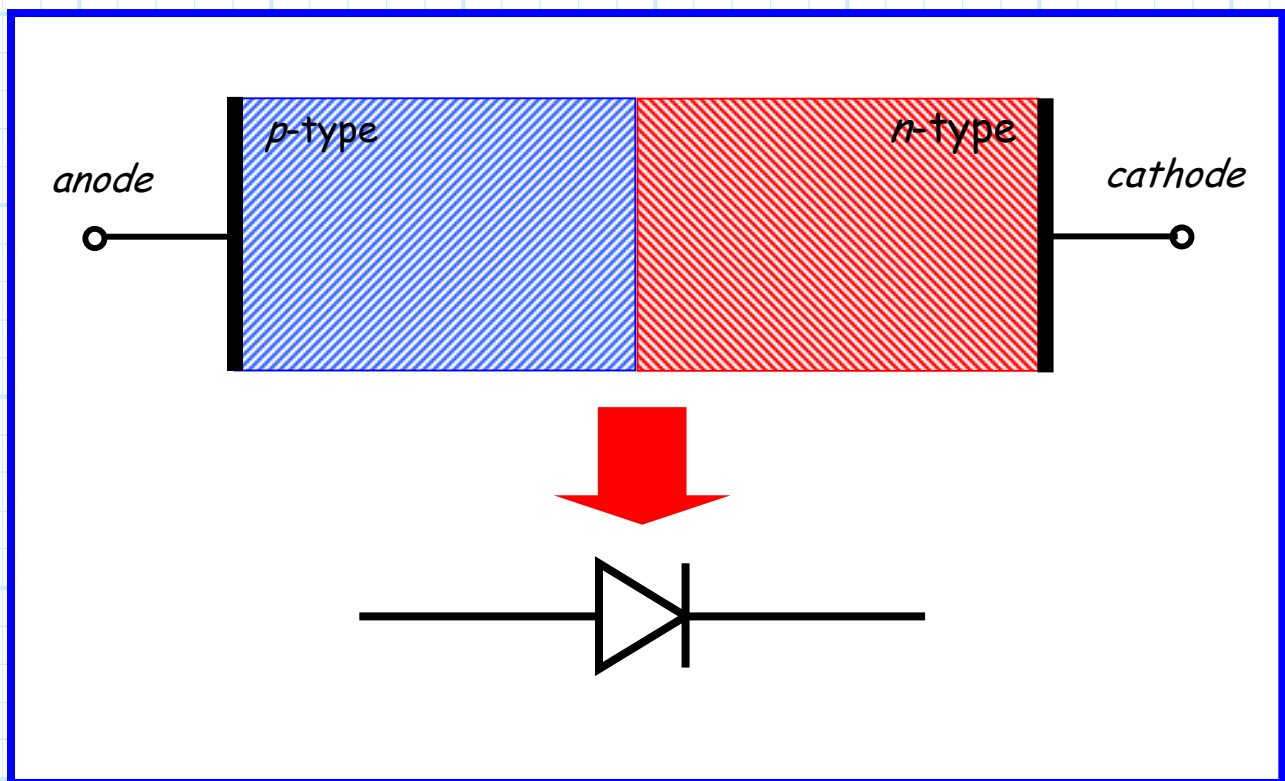


# The $p$ - $n$ Junction Diode (Open Circuit)

We create a  $p$ - $n$  junction diode simply by sticking together a hunk of  $p$ -type Silicon and a hunk of  $n$ -type Silicon!



Now, let's think about what happens here:

- 1) The **concentration** of holes in the anode is much greater than that of the cathode.
- 2) The **concentration** of free electrons in the cathode is much greater than that of the anode.

**Diffusion** is the result !

- 1) **Holes** begin to migrate across the junction from the **anode** to the **cathode**.
- 2) **Free electrons** begin to migrate across the junction from the **cathode** to the **anode**.

**Q:** *Oh, I see! This is **entropy** at work. This diffusion will occur until the concentration of holes and free electrons become **uniform** throughout the diode, right ?*



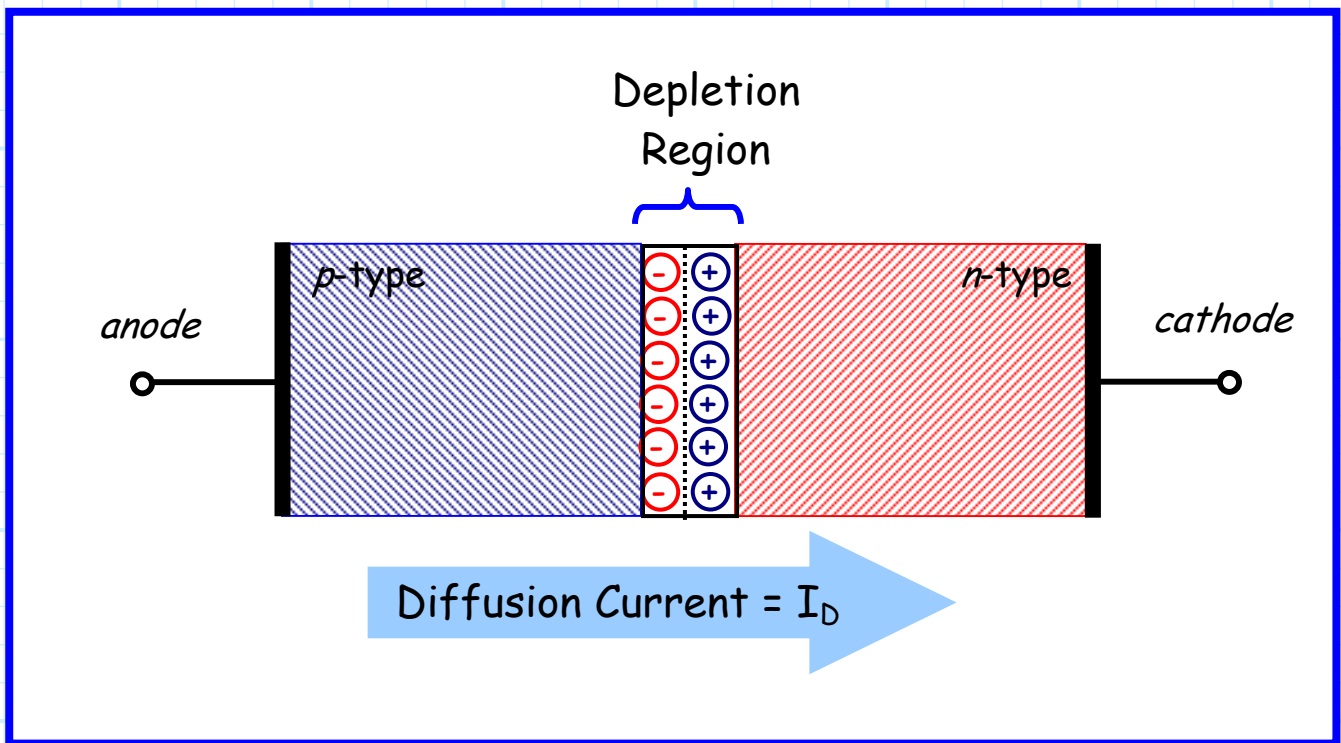
**A:** Not so fast ! There are **more** phenomena at work here than **just** diffusion !

For instance, **think** about what happens when holes **leave** the *p*-type Silicon of the anode, and the free electrons **leave** the *n*-type Silicon of the cathode:

**→** They **uncover ions !!!**

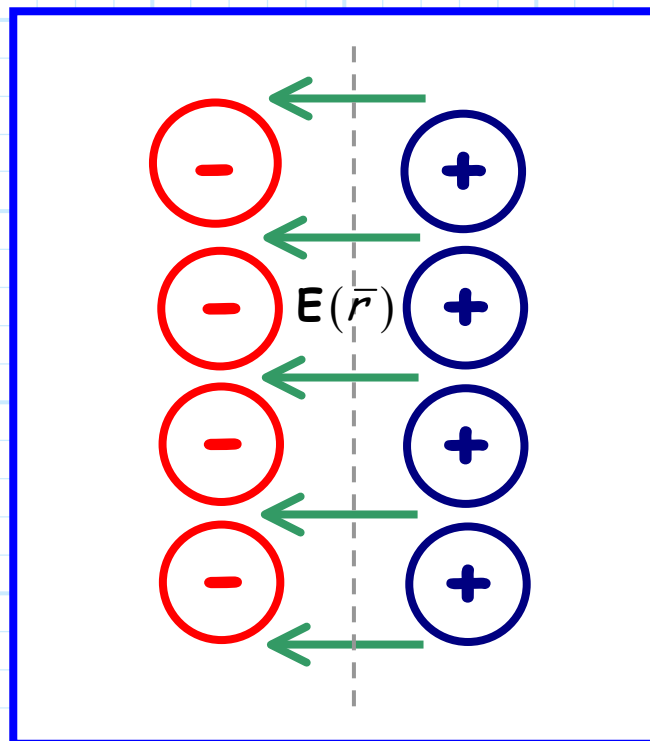
As a result, the **charge density** of the **anode** along the junction becomes **negative**, and the charge density of the **cathode** along the junction becomes **positive**.

This region of uncovered ions along the junction is known as the **depletion region**.



Now, something **really** interesting occurs.

The uncovered ions of opposite polarity generate an **electric field** across the junction !



Recall that an electric field exerts a **force** on charge particles—charged particles like **holes** and **free electrons**!

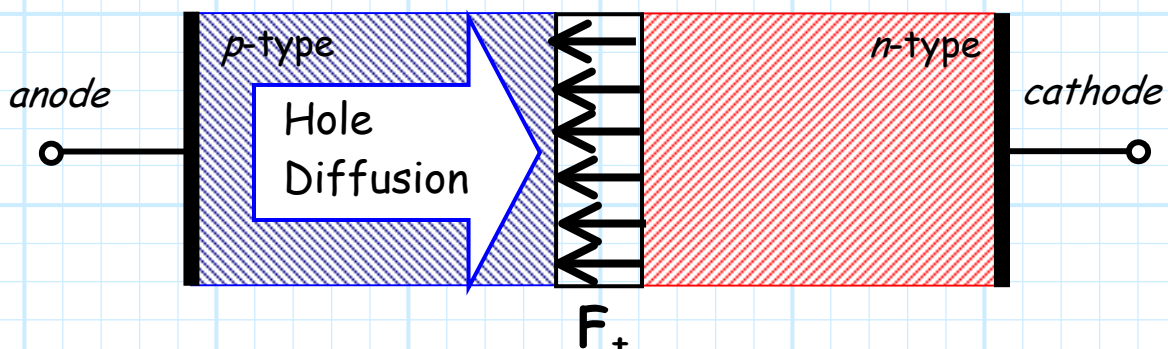
Let's see what this force is on both holes and free electrons:

### For holes:

Using the **Lorentz force equation**, we find that the force vector  $\mathbf{F}_+$  on a hole (with charge  $Q_+ = -e$ ) located at position  $\vec{r}$  is:

$$\mathbf{F}_+ = Q_+ \mathbf{E}(\vec{r})$$

Note that since the electric field vector in the depletion region is pointing from **right to left** (i.e., from the *n*-type Si cathode to the *p*-type Si anode), and since the "charge"  $Q_+$  of a hole is **positive**, the force vector likewise extends from **right to left**:



Look what happens! The electric field in the depletion region applies a **force** on the holes that is **opposite** of the direction of **hole diffusion**!

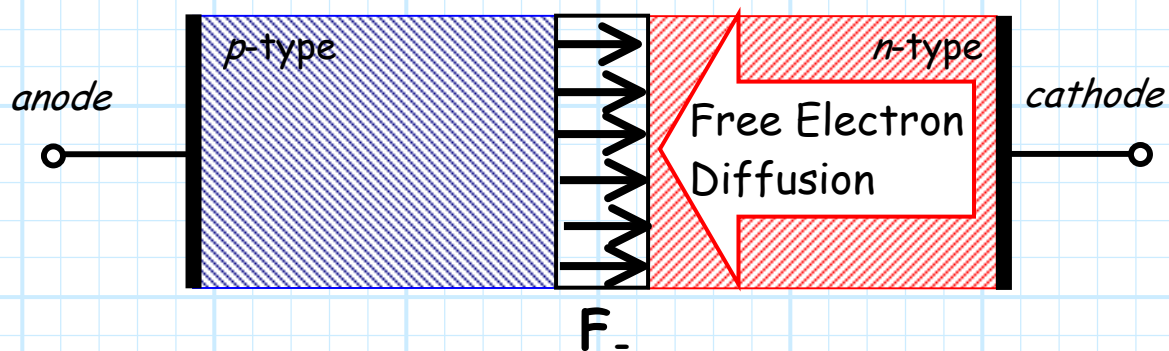
In other words, the electric field "**holds back**" the tide of holes attempting to **diffuse** into the *n*-type cathode region.

### For free electrons:

Now, let's see what effect this electric field has on **free electrons**. Using the **Lorentz force equation**, we find that the **force vector**  $\mathbf{F}_-$  on a free electron (with charge  $Q_- = e$ ) located at position  $\vec{r}$  is:

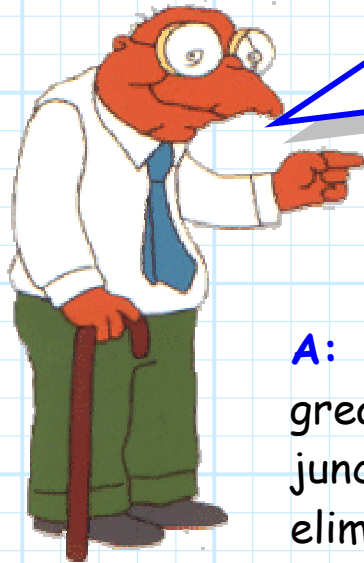
$$\mathbf{F}_- = Q_- \mathbf{E}(\vec{r})$$

Note that since the electric field vector in the depletion region is pointing from **right to left** (i.e., from the *n*-type Si to the *p*-type Si), and since the charge  $Q_-$  of a free electron is **negative**, the force vector extends in the opposite direction of  $\mathbf{E}(\vec{r})$ --from **left to right**:



Look what happens! The **electric field** in the depletion region likewise applies a **force** on the **free electrons** that is **opposite** of the direction of free electron diffusion!

In other words, the electric field "**holds back**" the tide of **free electrons** attempting to **diffuse** into the *p*-type anode region .



**Q:** So, does this electric field **stop** all diffusion across the junction? Is the diffusion current  $I_D$  therefore **zero**?

**A:** Typically **NO!** The electric field will greatly **reduce** the diffusion across the junction, but only in **certain** cases will it eliminate  $I_D$  entirely (more about **that** later!).

The **amount** of diffusion that occurs for a given electric field  $\mathbf{E}(\vec{r})$  is dependent on how **energetic** the particles (holes and free-electrons) are!

Recall that these particles will have **kinetic energy** due to heat. If this energy is sufficiently **large**, a particle can still diffuse **across** the  $p$ - $n$  junction!

To see why, consider the amount of **energy**  $E$  it would take to move a charged particle **through** this electric field. Recall from EECS 220 that this energy is:

$$E = -Q \int_C \mathbf{E}(\vec{r}) \cdot d\vec{\ell}$$

For our case,  $Q$  is the **charge** on a particle (hole or free electron), and **contour**  $C$  is a path that extends **across** the depletion region.

Moreover, we recall that this expression can be simplified by using **electric potential**, i.e.,

$$V = -\int_C \mathbf{E}(\bar{r}) \cdot \overline{d\ell}$$

Where  $V$  is the difference in **potential energy** (per coulomb) between a charge at either end of contour  $C$ . This of course tells us how much **work** must be done (per coulomb) to move a charge from **one end** of the contour to the **other**.

Of course  $V$  has units of **Volts**, but its more descriptive unit is **joules/coulomb**—energy per unit charge.

Therefore, the energy required to move a charge  $Q$  along some contour  $C$  can **likewise** be expressed as:

$$E = QV$$

Now, for our particular problem, the charge  $Q$  is either the charge of a **free electron** ( $Q_e$ ) or the charge of a **hole** ( $Q_h$ ).

The **voltage** (i.e., potential difference) across the depletion region is called the **barrier voltage**  $V_B$  (sometimes denoted as  $V_0$ ):

$$V_B = -\int_{C_{dr}} \mathbf{E}(\bar{r}) \cdot \overline{d\ell}$$

where the contour  $C_{dr}$  describes some path **across** the depletion region.

Typically, we find that when the junction diode is open circuited (i.e.,  $v_D=0$  and  $i_D=0$ ), this barrier voltage is approximately **−0.7 V** !

Thus, we find that the **energy** required for a **hole** to **diffuse** across the depletion region is:

$$E_B = Q_+ V_B$$

While the **energy** required for a **free electron** to **diffuse** across the depletion region is:

$$E_B = -Q_- V_B$$

Note that both these energies are the **same** (positive) value!

OK, here's the **important part**:

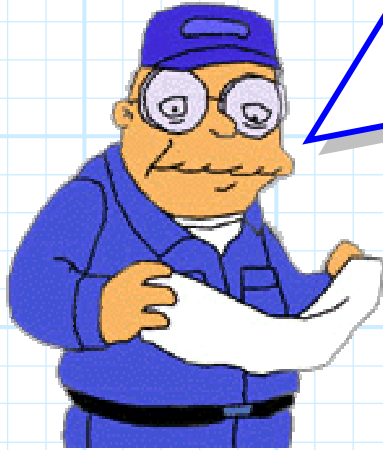
- A. **If** the particle has kinetic energy **greater** than  $E_B$ , it can **diffuse** across the depletion region.
- B. **If** the particle has kinetic energy **less** than  $E_B$ , then the electric field will "push" it **back** into either the  $p$ -type anode region (for holes) or the  $n$ -type cathode region (for free electrons).

Thus, the diffusion current  $I_D$  across the  $p$ - $n$  junction will depend on three things:

1. **The majority particle concentration.** - The more holes or free electrons there are, the more particles will diffuse across the junction.
2. **The barrier voltage  $V_B$ .** - A lower barrier means more diffusion, a higher means less.



3. **The diode temperature** - Higher temperature means more kinetic energy and thus more diffusion. Lower temperature means less.



**Q:** *Wait a minute! We've examined the behavior of holes in the p-type region and free electrons in the n-type region. These are the **majority** carriers for each of those Silicon types.*

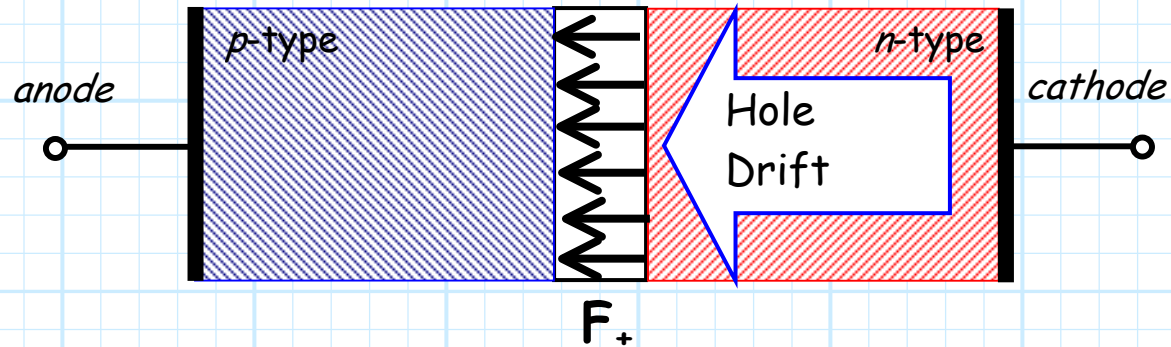
*There are also **minority** carriers present in each side. What does the electric field in the depletion region do to them?*

**A:** A great question! We will find that the electric field will have a profoundly **different** effect on **minority** carriers!

### For holes:

Recall that the electric field in the depletion region applies a force on **positive** charges (holes) that is directed **from** the *n*-type (cathode) region **into** the *p*-type (anode) region.

This force of course **pushes** the holes in the *p*-type anode (the *majority* carriers) **back into** the *p*-type region. **However**, the same force will **pull** holes from the *n*-type region (the **minority** carriers) **into** the *p*-type region!



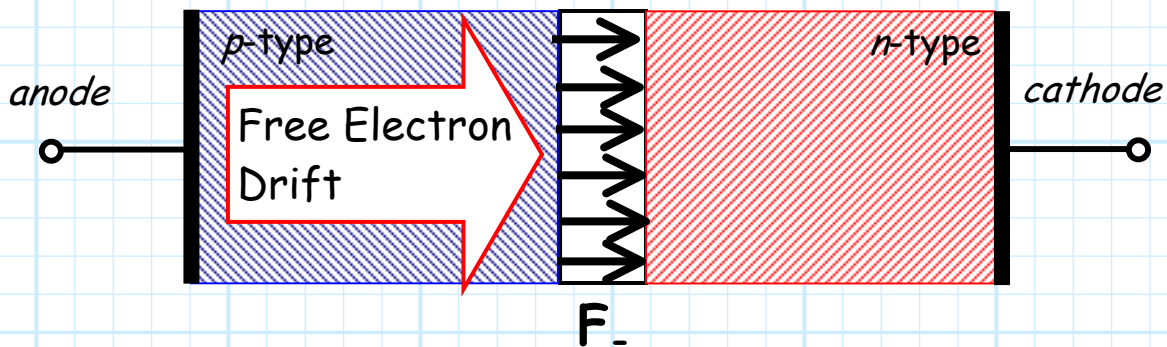
Any unsuspecting **minority** hole that "drifts" into the depletion region will from the *n*-type side will be **pulled** into the *p*-type side! Note that this result is **independent** of the kinetic energy of the particle—it takes **no** energy to "fall **downhill**".

This movement of charge is completely due to the force applied by the **electric field**—this is **drift current**  $I_s$ !

### Now, for free electrons:

Recall also that the electric field in the depletion region applies a force on **negative** charges (free electrons) that is directed **from** the *p*-type (anode) region **into** the *n*-type (cathode) region.

This force of course **pushes** the free electrons in the *n*-type region (the **majority** carriers) **back into** the *n*-type region. **However**, the same force will **pull** free electrons from the *p*-type region (the **minority** carriers) into the *n*-type region!

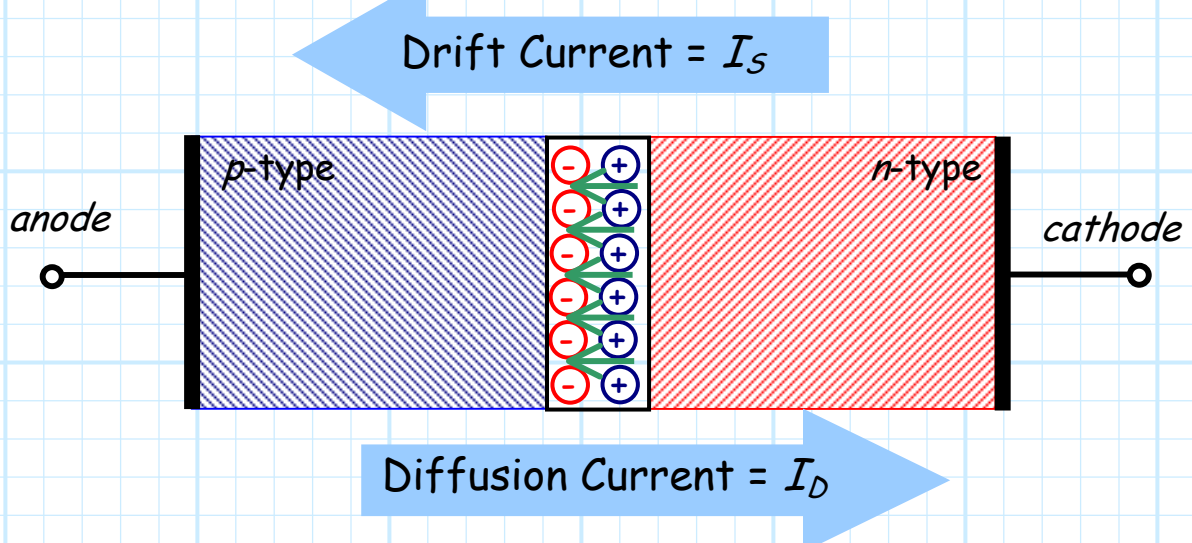


Any unsuspecting **minority** free electron that "drifts" into the depletion region will from the  $p$ -type side will be **pulled** into the  $n$ -type side! Note that this result is **independent** of the kinetic energy of the particle—it takes **no energy** to "fall downhill".

This movement of charge is completely due to the force applied by the **electric field**—this is also **drift current**  $I_S$ !

There are two very important **differences** between **drift** and **diffusion** currents in a  $p$ - $n$  junction diode:

- 1. Drift and Diffusion current flow in opposite directions** - The Diffusion current  $I_D$  flows across the  $p$ - $n$  junction from anode to cathode, while Drift current  $I_S$  flows across the  $p$ - $n$  junction from cathode to anode.



**2. Diffusion current depends on the barrier voltage  $V_B$ , but Drift Current does not.** - As the barrier voltage increases, fewer and fewer of the majority carriers will have sufficient kinetic energy to cross the depletion region—the **diffusion current will decrease.**

Conversely, minority carriers require **no energy** to be swept across the depletion region by the electric field, the value of the **barrier voltage is irrelevant** to the value of  $I_S$ .

Now, for an open-circuited (i.e., **disconnected**) junction diode, the **total current** through the device **must be zero**. In other words, the diffusion current  $I_D$  must be **equal but opposite** that of the drift current  $I_S$ , such that:

$$I_D - I_S = 0$$

This is the **equilibrium** state of a **disconnected** junction diode. We find that typically this drift/diffusion current is **very small**, generally  $10^{-8}$  to  $10^{-12}$  Amps!