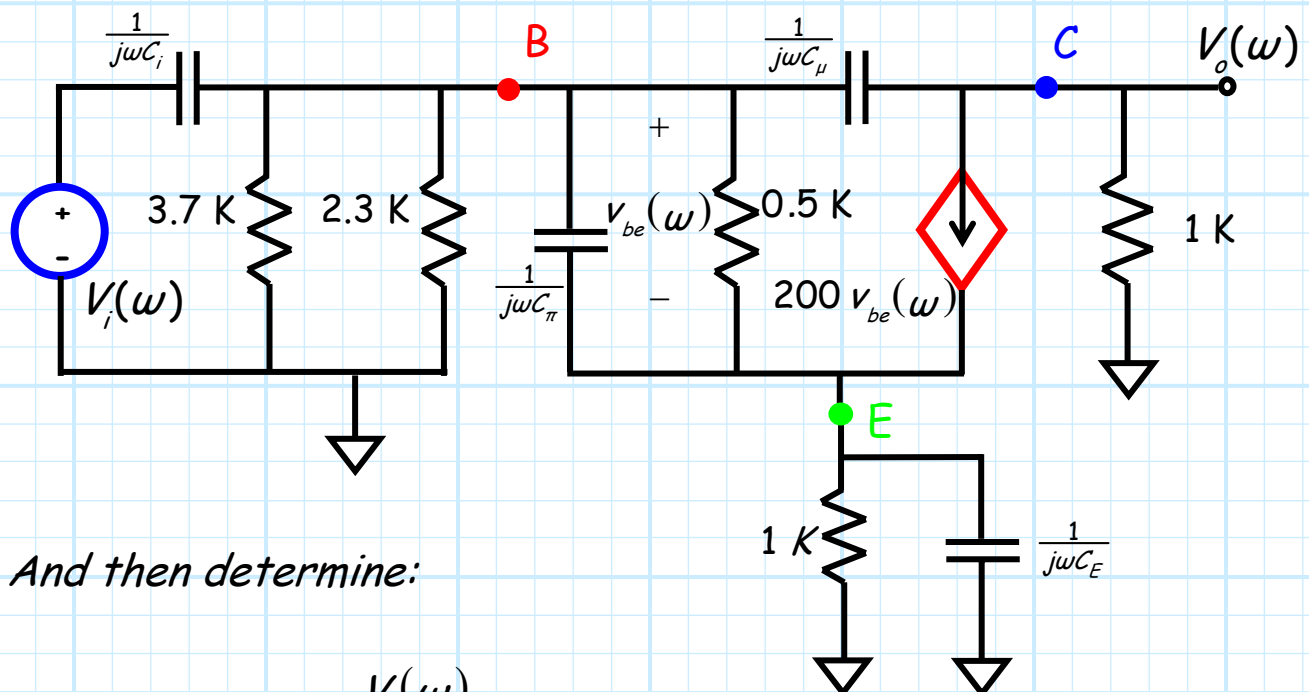


Low-Frequency Response

Q: OK, I see how to determine mid-band gain, but what about determining amplifier **bandwidth**?

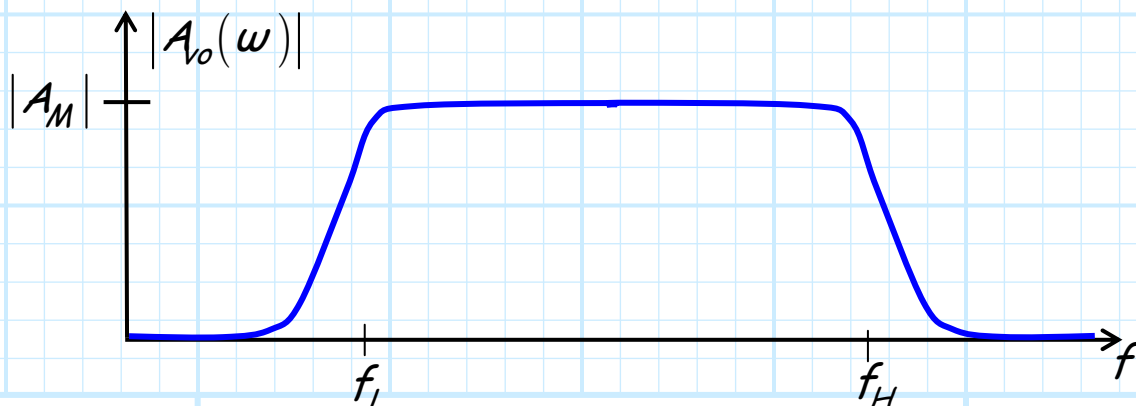
It seems like I have no alternative but to analyze the exact small-signal circuit (explicitly considering all capacitances):



And then determine:

$$A_{vo}(\omega) = \frac{V_o(\omega)}{V_i(\omega)}$$

and then **plot** the magnitude:



And then from the plot determine the amplifier bandwidth (i.e., determine f_L and f_H)?

A: You could do all that, but there is an easier way.

An amplifier frequency response $A_o(\omega)$ (i.e., its eigen value!) can generally be expressed as the product of **three** distinct terms:

$$A_o(\omega) = F_L(\omega) A_M F_H(\omega)$$

The middle term is of course the **mid-band gain**—a number that is not frequency dependent.

The function $F_L(\omega)$ describes the **low-frequency response** of the amplifier—from it we can determine the lower cutoff frequency f_L .

Conversely, the function $F_H(\omega)$ describes the **high-frequency response** of the amplifier—from it we can determine the upper cutoff frequency f_H .

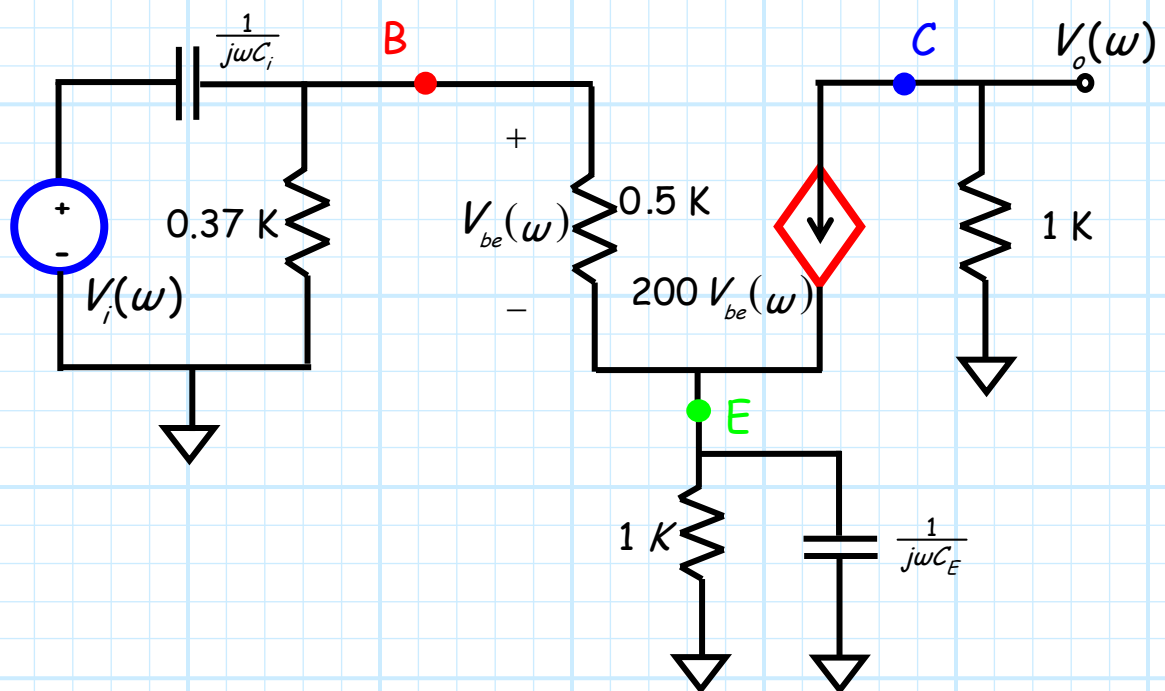
Q: *So just how do we determine these functions $F_L(\omega)$ and $F_H(\omega)$??*

A: The low-frequency response $F_L(\omega)$ is dependent **only** on the large capacitors (COUS) in the amplifier circuit. In other

words the parasitic capacitances have **no affect** on the low-frequency response.

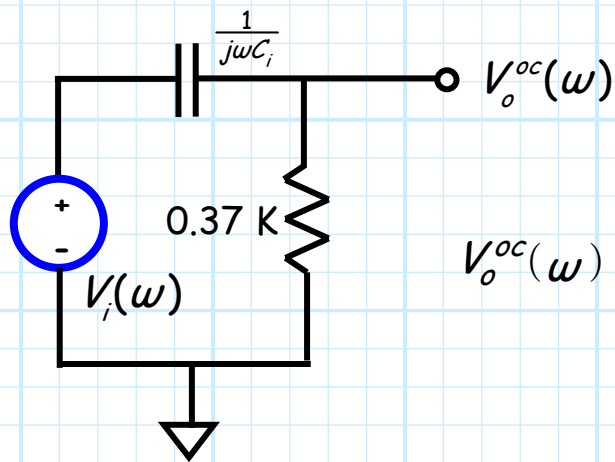
Thus, we simply "ignore" the **parasitic** capacitances when determining $F_L(\omega)$!

For example, say we include the COUS in our common-emitter example, but **ignore** C_μ and C_π . The resulting small-signal circuit is:



To simplify this analysis, we first determine the **Thevenin's** equivalent circuit of the portion of the circuit connected to the base.

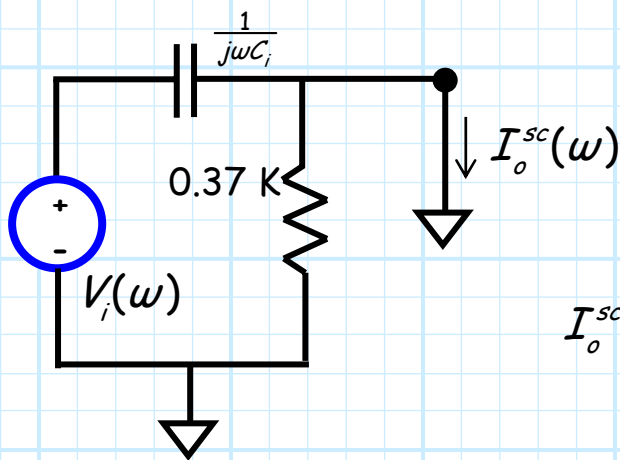
We start by finding the **open-circuit voltage**:



$$V_o^{oc}(\omega) = V_i(\omega) \left(\frac{0.37}{0.37 + \frac{1}{j\omega C_i}} \right)$$

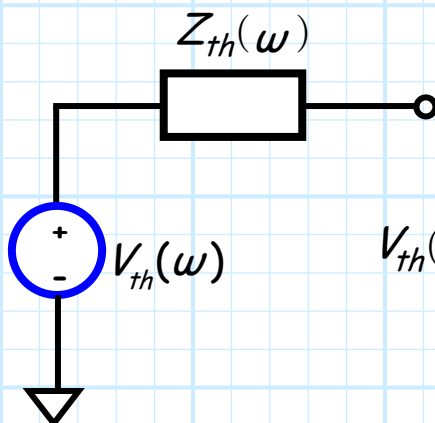
$$= V_i(\omega) \left(\frac{j\omega(0.37)C_i}{1 + j\omega(0.37)C_i} \right)$$

And the short-circuit output current is:



$$I_o^{sc}(\omega) = \frac{V_i(\omega)}{\frac{1}{j\omega C_i}} = V_i(\omega) (j\omega C_i)$$

And thus the Thevenin's equivalent source is:



$$V_{th}(\omega) = V_o^{oc}(\omega)$$

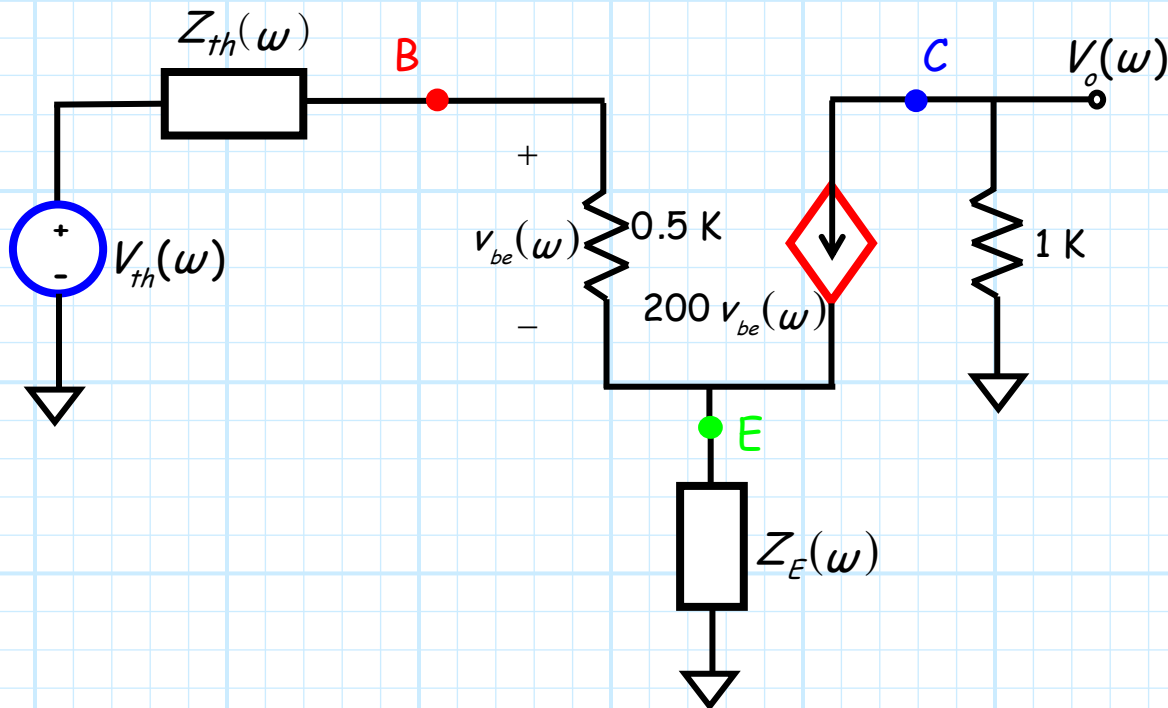
$$= V_i(\omega) \left(\frac{j\omega(0.37)C_i}{1 + j\omega(0.37)C_i} \right)$$

$$\begin{aligned}
 Z_{th}(\omega) &= \frac{V_o^{oc}(\omega)}{I_o^{sc}(\omega)} \\
 &= \left(\frac{j\omega(0.37)C_i}{1 + j\omega(0.37)C_i} \right) \left(\frac{1}{j\omega C_i} \right) \\
 &= \frac{(0.37)}{1 + j\omega(0.37)C_i}
 \end{aligned}$$

Likewise, the two parallel elements on the emitter terminal can be combined:

$$Z_E(\omega) = \frac{\frac{1}{j\omega C_E}}{1 + \frac{1}{j\omega C_E}} = \frac{1}{1 + j\omega C_E}$$

Thus, the small-signal circuit is now:



From KVL:

$$0 + V_{th} - I_b (Z_{th} + 0.5) - (\beta + 1) I_b Z_E = 0$$

$$\Rightarrow I_b = \frac{V_{th}}{Z_{th} + 0.5 + 101Z_E}$$

From Ohm's Law:

$$V_{be} = 0.5I_b$$

Therefore:

$$\begin{aligned} V_o(\omega) &= -200 V_{be}(1) \\ &= -200(0.5) \frac{V_{th}(\omega)}{Z_{th} + 0.5 + 101Z_E} \\ &= V_{th}(\omega) \left(\frac{-100}{Z_{th} + 0.5 + 101Z_E} \right) \end{aligned}$$

Inserting the expressions for the Thevenin's equivalent source, as well as Z_E .

$$\begin{aligned} V_o(\omega) &= V_{th}(\omega) \left(\frac{-100}{Z_{th} + 0.5 + 101Z_E} \right) \\ &= V_i(\omega) \left(\frac{j\omega(0.37)C_i}{1 + j\omega(0.37)C_i} \right) \left(\frac{-200}{\frac{2(0.37)}{1 + j\omega(0.37)C_i} + 1 + \frac{202}{1 + j\omega C_E}} \right) \end{aligned}$$

Now, it can be shown that:



$$\frac{1}{\frac{2(0.37)}{1 + j\omega(0.37)C_i} + 1 + \frac{202}{1 + j\omega C_E}} \approx \frac{j\omega(C_E/203)}{1 + j\omega(C_E/203)}$$

Therefore:

$$V_o(\omega) = V_i(\omega) \left(\frac{j\omega(0.37)C_i}{1 + j\omega(0.37)C_i} \right) \left(\frac{j\omega(C_E/203)}{1 + j\omega(C_E/203)} \right) (-200)$$

And so:

$$A_{v_o}(\omega) = \left(\frac{j\omega(0.37)C_i}{1 + j\omega(0.37)C_i} \right) \left(\frac{j\omega(C_E/203)}{1 + j\omega(C_E/203)} \right) (-200)$$

Now, since we are **ignoring** the parasitic capacitances, the function $F_H(\omega)$ that describes the high frequency response is:

$$F_H(\omega) = 1$$

And so:

$$A_{v_o}(\omega) = F_L(\omega) A_M F_H(\omega) = F_L(\omega) A_M$$

By inspection, we see for this example:

$$A_M = -200 \quad \leftarrow \text{We knew this already!}$$

And:

$$F_L(\omega) = \left(\frac{j\omega(0.37)C_i}{1 + j\omega(0.37)C_i} \right) \left(\frac{j\omega(C_E/203)}{1 + j\omega(C_E/203)} \right)$$

Now, let's define:

$$\omega_{P1} = \frac{1}{0.37C_i} = \frac{2.7}{C_i} \quad \text{and} \quad \omega_{P2} = \frac{203}{C_E}$$

Thus,

$$F_L(\omega) = \left(\frac{j(\omega/\omega_{P1})}{1 + j(\omega/\omega_{P1})} \right) \left(\frac{j(\omega/\omega_{P2})}{1 + j(\omega/\omega_{P2})} \right)$$

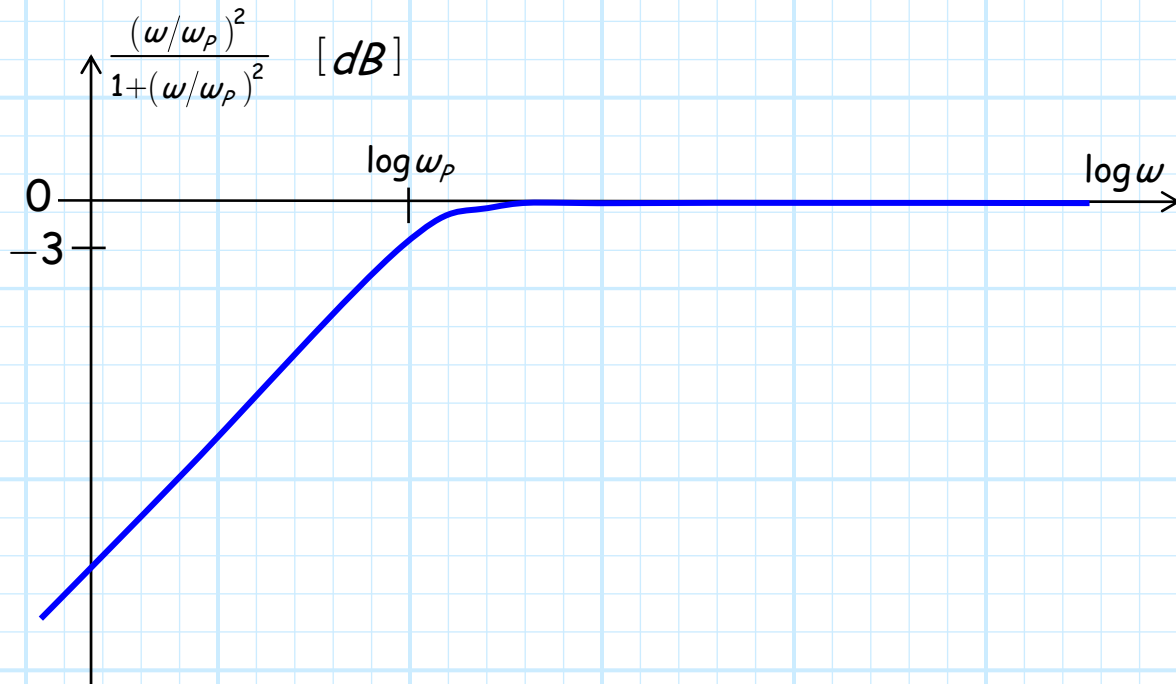
Now, functions of the type:

$$\left(\frac{j(\omega/\omega_p)}{1 + j(\omega/\omega_p)} \right)$$

are **high-pass** functions:

$$\left| \frac{j(\omega/\omega_p)}{1 + j(\omega/\omega_p)} \right|^2 = \frac{(\omega/\omega_p)^2}{1 + (\omega/\omega_p)^2}$$

with a **3dB break frequency** of ω_p .



Thus:

$$\frac{(\omega/\omega_p)^2}{1 + (\omega/\omega_p)^2} = \begin{cases} \approx 1.0 & \text{for } \omega > \omega_p \\ 0.5 & \text{for } \omega = \omega_p \\ 0 & \text{for } \omega = 0 \end{cases}$$

As a result, we find that the transfer function:

$$\begin{aligned} A_v(\omega) &= F_L(\omega) A_M \\ &= \left(\frac{j(\omega/\omega_{p1})}{1 + j(\omega/\omega_{p1})} \right) \left(\frac{j(\omega/\omega_{p2})}{1 + j(\omega/\omega_{p2})} \right) (-200) \end{aligned}$$

will be **approximately** equal to the midband gain $A_M = -200$ for all frequencies ω that are greater than **both** ω_{p1} and ω_{p2} .

I.E.,:

$$A_{vo}(\omega) \cong A_M = -200 \quad \text{if } \omega > \omega_{p1} \text{ and } \omega > \omega_{p2}$$

Hopefully, it is **now** apparent (please tell me it is!) that the lower end of the amplifier bandwidth—specified by frequency ω_L —is determined by the **larger** of the two frequencies ω_{p1} and ω_{p2} !

The **larger** of the two frequencies is called the **dominant pole** of the transfer function $F_L(\omega)$.

For our example—comparing the two frequencies ω_{p1} and ω_{p2} :

$$\omega_{p1} = \frac{1}{0.37C_i} = \frac{2.7}{C_i} \quad \text{and} \quad \omega_{p2} = \frac{203}{C_E}$$

it is apparent that the **larger** of the two (the dominant pole!) is likely ω_{p2} —that darn **emitter capacitor** is the key!

Say we want the common-emitter amplifier in this circuit to have a bandwidth that extends down to $f_L = 100 \text{ Hz}$

The emitter capacitor **must** therefore be:

$$2\pi f_L > \omega_{p2} = \frac{203}{C_E}$$

$$\Rightarrow C_E > \frac{203}{2\pi f_L} = \frac{203}{2\pi(1000)} = 32,300\mu F \quad !!!$$

This certainly is a **Capacitor Of Unusual Size** !

