### B. Amplifiers

We will find that the signal power collected by a receiver antenna is often **ridiculously small** (e.g., less than one **trillionth** of a Watt!)

To accurately recover the information impressed on this signal, we must **increase** the signal power a whole bunch—**without** modifying or distorting the signal in any way.

We accomplish this with a RF/microwave amplifier—one of the few **active** components we will study.

But first, a few comments about the **decibel**!

- HO: dB, dBm, dBw
- HO: Amplifiers
- Q: By how much will an amplifier increase signal power?
- A: <u>HO: Amplifier Gain</u>
- Q: Can we increase this signal power an unlimited amount?

A: NO! At some point we are limited by conservation of energy!

### HO: Amplifier Output Power

**Q:** So, just how precisely does an amplifier reproduce a signal at its output?

A: HO: Intermodulation Distortion

Q: Is intermodulation distortion really that big of a problem?

A: It can be if there are **multiple** signals at the amplifier input!

HO: Two-Tone Intermodulation Distortion

Every good radio engineer knows and understands that parameters of the amplifier **spec sheet**!

HO: The Amplifier Spec Sheet

# dB. dBm. dBw

**Decibel** (dB), is a specific function that operates on a **unitless** parameter:

$$dB \doteq 10 \ \log_{10}(x)$$

where x is <u>unitless</u>!

Q: A unitless parameter! What good is that !?

A: Many values are unitless, such as ratios and coefficients.

For example, amplifier gain is a unitless value!

E.G., amplifier gain is the ratio of the output power to the input power:

$$\frac{P_{out}}{P_{in}} = G$$

$$\therefore$$
 Gain in dB = 10  $\log_{10}\mathcal{G} \doteq \mathcal{G}(dB)$ 

Jim Stiles

**Q:** Wait a minute! I've seen statements such as:

.... the output **power** is 5 dBw ....

or .... the input **power** is 17 dBm ....

Of course, Power is **not** a unitless parameter!?!

A: True! But look at how power is expressed; not in dB, but in dBm or dBw.

Q: What the heck does dB<u>m</u> or dB<u>w</u> refer to ??

A: It's sort of a trick !

Say we have some power *P*. Now say we divide this value *P* by one **1 Watt**. The result is a **unitless** value that expresses the value of *P* in **relation** to 1.0 Watt of power.

For **example**, if  $P = 2500 \, mW$ , then P/1W = 2.5. This simply means that power P is 2.5 times larger than one Watt!

Since the value P/1W is **unitless**, we can express this value in **decibels**!

Specifically, we define this operation as:

$$P(dBw) \doteq 10 \log_{10}\left(\frac{P}{1W}\right)$$

For example, P = 100 Watts can alternatively be expressed as P(dBw) = +20 dBw. Likewise, P = 1 mW can be expressed as P(dBw) = -30 dBw.

### Q: OK, so what does dBm mean?

A: This notation simply means that we have normalized some power P to one Milliwatt (i.e., P/1mW)—as opposed to one Watt. Therefore:

$$P(dBm) \doteq 10 \log_{10}\left(\frac{P}{1 \ mW}\right)$$

For example, P = 100 Watts can alternatively be expressed as P(dBm) = +50 dBm. Likewise, P = 1 mW can be expressed as P(dBm) = 0 dBm.

Make sure you are very **careful** when doing math with decibels!

### Standard dB Values

Note that  $10 \log_{10} (10) = 10 dB$ 

Therefore an amplifier with a gain G = 10 is likewise said to have a gain of 10 dB.

Now consider an amplifier with a gain of **20** dB.....



**Q:** Yes, yes, I know. A 20 dB amplifier has gain G=20, a 30 dB amp has G=30, and so forth.

Please speed this lecture up and quit wasting my valuable time making such **obvious** statements!



A: NO! Do not make this mistake!



Recall from your knowledge of logarithms that:

 $10\log_{10}[10^n] = n \, 10\log_{10}[10] = 10n$ 

Therefore, if we express gain as  $G = 10^n$ , we conclude:

$$\mathcal{G} = 10^n \quad \leftrightarrow \quad \mathcal{G}(dB) = 10n$$

In other words,  $G=100 = 10^2$  (n=2) is expressed as 20 dB, while 30 dB (n=3) indicates  $G=1000 = 10^3$ .

Likewise 100 mW is denoted as 20 dBm, and 1000 Watts is denoted as 30 dBW.

Note also that 0.001 mW =  $10^{-3}$  mW is denoted as -30 dBm.

Another important relationship to keep in mind when using decibels is  $10 \log_{10}{[2]} \approx 3.0$  . This means that:

 $10\log_{10}[2^n] = n \ 10\log_{10}[2] \simeq 3n$ 

Therefore, if we express gain as  $G = 2^n$ , we conclude:

$$\mathcal{G} = 2^n \quad \leftrightarrow \quad \mathcal{G}(dB) \simeq 3n$$

As a result, a **15 dB** (n=5) gain amplifier has  $G = 2^5 = 32$ . Similarly, **1/8** =  $2^{-3}$  mW (n=-3) is denoted as -**9 dBm**.

Jim Stiles

### **Multiplicative Products and Decibels**

Other logarithmic relationships that we will find useful are:

$$10\log_{10}[xy] = 10\log_{10}[x] + 10\log_{10}[y]$$

and its close cousin:

$$10\log_{10}\left|\frac{x}{y}\right| = 10\log_{10}[x] - 10\log_{10}[y]$$

Thus, the relationship  $P_{out} = G P_{in}$  is written in **decibels** as:

$$P_{out} = \mathcal{G} P_{in}$$

$$\frac{P_{out}}{1mW} = \frac{\mathcal{G} P_{in}}{1mW}$$

$$10\log_{10} \left[ \frac{P_{out}}{1mW} \right] = 10\log_{10} \left[ \frac{\mathcal{G} P_{in}}{1mW} \right]$$

$$10\log_{10} \left[ \frac{P_{out}}{1mW} \right] = 10\log_{10} \left[ \mathcal{G} \right] + 10\log_{10} \left[ \frac{P_{in}}{1mW} \right]$$

$$P_{out}(dBm) = G(dB) + P_{out}(dBm)$$

It is evident that "deebees" are **not** a unit! The units of the result can be found by **multiplying** the units of each term in a **summation** of decibel values.

For example, say some power $P_1 = 6  dBm$ is <b>combined</b> with
power $P_2 = 10  dBm$ . What is the resulting <b>total</b> power $P_T = P_1 + P_2$ ?
O: This result really is obvious-
of course the total power is:
$P_{T}(dBm) = P_{1}(dBm) + P_{2}(dBm)$
= 6  dBm + 10  dBm
= 16  dBm
A: NO! Never do this either!
Logarithms are very helpful in expressing <b>products</b> or <b>ratios</b>
involves sums and differences!
$10\log_{10}[x+y] = ????$
So, if you wish to add $P_1$ =6 dBm of power to $P_2$ =10 dBm of power, you must first <b>explicitly</b> express power in Watts:
$P_1 = 10 \ dBm = 10 \ mW$ and $P_2 = 6 \ dBm = 4 \ mW$

Thus, the total power  $P_T$  is:

$$P_T = P_1 + P_2$$
  
= 4.0 mW + 10.0 mW  
= 14.0 mW

Now, we can express this total power in *dBm*, where we find:

$$P_{T}(dBm) = 10 \log_{10}\left(\frac{14.0 \ mW}{1.0 \ mW}\right) = 11.46 \ dBm$$

The result is **not** 16.0 *dBm* !.

We **can** mathematically add 6 *dBm* and 10 *dBm*, but we must understand what result means (**nothing useful**!).

$$6 dBm + 10 dBm = 10 \log_{10} \left[ \frac{4mW}{1mW} \right] + 10 \log_{10} \left[ \frac{10mW}{1mW} \right]$$
$$= 10 \log_{10} \left[ \frac{40 mW^2}{1mW^2} \right]$$

= 16 *dB* relative to 1 mW<sup>2</sup>

Thus, mathematically speaking, 6 *dBm* + 10 *dBm* implies a multiplication of power, resulting in a value with units of **Watts squared**!

9/5/2007



# <u>Amplifiers</u>

An **ideal** amplifier takes an input signal and reproduces it **exactly** at its output, only with a **larger** magnitude!





where  $A_v$  is the voltage gain of the amplifier.

Now, let's express this result using our knowledge of linear circuit theory !

Recall, the output  $v_{out}(t)$  of a linear device can be determined by convolving its input  $v_{in}(t)$  with the device impulse response h(t):

$$\mathbf{v}_{out}(t) = \int h(t-t')\mathbf{v}_{in}(t')dt'$$

The impulse response for the **ideal** amplifier would therefore be:  $h(t) = A \delta(t)$ so that:



→ Any and every function  $v_{in}(t)$  is an eigen function of an ideal amplifier!!

We can alternatively represent the ideal amplifier response in the **frequency domain**, by taking the **Fourier Transform** of the impulse response:

$$\mathcal{H}(\omega) = \int_{-\infty}^{\infty} h(t) e^{-j\omega t} dt$$
$$= \int_{-\infty}^{\infty} \mathcal{A} \delta(t) e^{-j\omega t} dt$$
$$= \mathcal{A} + j0$$

This result, although simple, has an interesting interpretation. It means that the amplifier exhibits gain of  $A_{\nu}$  for sinusoidal signals of **any** and **all** frequencies!

$$\left|\mathcal{H}(\omega)\right|^{\uparrow}$$

BUT, there is one **big** problem with an ideal amplifier:

They are **impossible** to build !!

The **ideal** amplifier has a frequency response of  $|\mathcal{H}(\omega)| = \mathcal{A}_{\nu}$ . Note this means that the amplifier gain is  $\mathcal{A}_{\nu}$  for **all** frequencies  $0 < \omega < \infty$  (D.C. to daylight !).

The bandwidth of the ideal amplifier is therefore infinite !

\* Since every electronic device will exhibit **some** amount of inductance, capacitance, and resistance, every device will have a **finite** bandwidth.

\* Moreover, we discussed that a matching network likewise exhibits **finite** bandwidth.

\* In other words, there will be frequencies ω where the device does not work!

\* From the standpoint of an amplifier, "not working" means  $|\mathcal{H}(\omega)| \ll A_{\mu}$  (i.e., **low gain**).

> Amplifiers therefore have **finite** bandwidths.

There is a range of frequencies  $\omega$  between  $\omega_L$  and  $\omega_H$  where the gain will (approximately) be  $A_{\nu}$ . For frequencies outside this range, the gain will typically be small (i.e.  $|\mathcal{H}(\omega)| \ll A_{\nu}$ ):

$$|\mathcal{H}(\omega)| = \begin{cases} \approx \mathcal{A}, & \omega_L < \omega < \omega_H \\ \ll \mathcal{A}, & \omega < \omega_L, \infty > \omega_H \end{cases}$$
  
The width of this frequency range is called the amplifier bandwidth:  
Bandwidth:  
Bandwidth  $\doteq \omega_H - \omega_L$  (radians/sec)  
 $\pm f_L - f_H$  (cycles/sec)  
 $|\mathcal{H}(\omega)| \uparrow$   
 $\mathcal{A}_{o} -$   
 $\mathcal{O}_L$   $\mathcal{O}_H$   
One result of having a finite bandwidth is that the amplifier impulse response is not an impulse function !

$$h(t) = \int \mathcal{H}(\omega) \, e^{+j\omega t} dt \neq \mathcal{A}_{\omega} \, \delta(t)$$

--∞

therefore generally speaking:

$$\mathbf{v}_{out}(t) \neq \mathbf{A}_{v} \mathbf{v}_{in}(t)$$

The ideal amplifier is not possible!

# Amplifier Gain

Note that an amplifier is a **two-port** device.



As a result, we can describe an amplifier with a 2 x 2 scattering matrix:

$$\boldsymbol{S} = \begin{bmatrix} \boldsymbol{S}_{11} & \boldsymbol{S}_{12} \\ \boldsymbol{S}_{21} & \boldsymbol{S}_{22} \end{bmatrix}$$

Q: What is the scattering matrix of an ideal amplifier??

A: Let's start with  $S_{11}$  and  $S_{22}$ .

To insure maximum power transfer, the input and output ports would ideally be matched:

$$S_{11} = S_{22} = 0$$

Now, let's look at scattering parameter  $S_{21}$ . We know that **if** the amplifier is connected to matched devices:

 $P_2^- = |S_{21}|^2 P_1^+$ 

or, stated **another** way:

$$P_{out} = \left| S_{21} \right|^2 P_{in}$$

Therefore, we can **define** the amplifier **power gain** as:

$$\boldsymbol{\mathcal{G}} \doteq \frac{\boldsymbol{\mathcal{P}}_{out}}{\boldsymbol{\mathcal{P}}_{in}} = \left|\boldsymbol{\mathcal{S}}_{21}\right|^2$$

As the purpose of an amplifier is to boost the signal power, we can conclude that **ideally**:

$$|\mathcal{S}_{21}| \gg 1$$

Clearly, an amplifier must be an **active** device!

As discussed earlier, the gain of an amplifier will change with signal frequency:

$$\boldsymbol{\mathcal{G}}(\boldsymbol{\omega}) = \left|\boldsymbol{\mathcal{S}}_{21}(\boldsymbol{\omega})\right|^2$$

When radio engineers speak of amplifier **gain**, they almost always are speaking of this **power gain G**. However, they do not generally state it as a specific function of frequency!

Rather, amplifier gain is typically specified as a **numeric** value such as G = 20 or G = 13 dB. This value is a statement of the approximate amplifier gain **within** the amplifier **bandwidth**.

G

 $\mathbf{G}(\omega)$ 

Thus, amplifier **gain** and **bandwidth** are the two most fundamental performance specifications of any microwave amplifier—together they (approximately) describe the amplifier transfer function!

 $\omega_{l}$ 

 $\dot{\omega}_{H}$ 

Additionally, radio engineers almost always speak of amplifier gain in **decibels** (dB):

$$\mathcal{G}(dB) = 10 \ \log_{10} \mathcal{G}$$

Finally, let's consider  $S_{12}$ . This scattering parameter relates the wave into port 2 (the output) to the wave out of port 1 (the input).

 $V_{2}^{+}$ 



Ф

A: No! An amplifier is strictly a directional device; there is a specific input, and a specific output—it does not work in reverse!

**Ideally**,  $S_{12} = 0$ . Any other value can just cause problems!

**Typically** though,  $S_{12}$  is small, but **not** zero. Generally speaking, radio engineers express  $S_{12}$  as a value called **reverse isolation**:

reverse isolation  $\doteq -10 \log_{10} |S_{12}|^2$ 

Note when  $S_{12} = 0$ , reverse isolation will be **infinite**. Thus, the **larger** the reverse isolation, the **better**!

Summarizing, we find that the scattering matrix of the ideal amplifier is:

$$\boldsymbol{\mathcal{S}}_{ideal} = \begin{bmatrix} 0 & 0 \\ \boldsymbol{\mathcal{S}}_{21} & 0 \end{bmatrix}$$
 where  $|\boldsymbol{\mathcal{S}}_{21}| \gg 1$ 

The non-ideal reality is that the zero valued terms will be small, but not precisely zero. Moreover, each scattering parameter will change with signal frequency—although they remain approximately constant within the amplifier bandwidth.

# **Amplifier Output Power**

Say we have an amplifier with gain G = 30 dB (i.e., G = 1000).

If the input power to this amplifier is 0 dBw (i.e.,  $P_{in}$  = 1W), then the output power is:

$$P_{in} G = P_{out}$$
  
(1 W) 1000 = 1000 W

Or, in dB:

$$P_{in} = \frac{1W}{G}$$

WOW! We created 999 Watts !

The energy crisis is solved !

Of course, the amplifier cannot create energy.

Q: Then, where does the power come from ???

A: The D.C. power supply ! (Every amplifier has one).





The 1 dB compression point is generally considered to be the **maximum power output** of the amplifier.

The input power at the 1 dB compression point is said to be the **maximum input power**  $(P_{in}^{\max})$  of the amplifier. We of course **can** put more than  $P_{in}^{\max}$  into the amplifier—but we **won't** get much more power out!



Note the equation  $P_{out}(dB) = [P_{in}(dB) + G(dB)] - 1 dB$  alone is **not sufficient** to determine the 1 dB compression point, as we have two uknowns ( $P_{in}$  and  $P_{out}$ ). We need **another** equation!

This second "equation" is the actual curve or table of data relating  $P_{in}$  to  $P_{out}$  for a specific amplifier.

### **Amplifier Efficiency**

We can define **amplifier efficiency** e as the ratio of the maximum output power ( $P_{1dB}$ ) to the D.C. power:

$$e = \frac{P_{1dB}}{P_{DC}}$$
 (don't use decibels here!)

For example, if e=0.4, then up to 40% of the D.C. power **can** be converted to **output power**, while the remaining 60% is converted to **heat**.

We require high power amps to be very efficient!

## **Intermodulation** Distortion

The 1 dB compression curve shows that amplifiers are only **approximately** linear.

Actually, this should be **obvious**, as amplifiers are constructed with transistors—**non-linear** devices!

So, instead of the ideal case:

$$V_{out} = A_{v} V_{in}$$

Actual amplifier behavior requires more terms to describe!

$$\boldsymbol{v}_{out} = \boldsymbol{A}_{v} \boldsymbol{v}_{in} + \boldsymbol{B} \boldsymbol{v}_{in}^{2} + \boldsymbol{C} \boldsymbol{v}_{in}^{3} + \cdots$$

This representation is simply a **Taylor Series** representation of the **non-linear** function:

$$V_{out} = f(V_{in})$$

**Q:** Non-linear! But I thought an amplifier was a linear device? After all, we characterized it with a scattering matrix! A: Generally speaking, the constants B, C, D, etc. are **very** small compared to the voltage gain  $A_v$ . Therefore, **if**  $v_{in}$  is likewise small, we can **truncate** the Taylor Series and **approximate** amplifier behavior as the linear function:

 $V_{out} \approx A_{V} V_{in}$ 

**BUT**, as  $v_{in}$  gets large, the values  $v_{in}^2$  and  $v_{in}^3$  will get **really** large! In that case, the terms  $B v_{in}^2$  and  $C v_{in}^3$  will become **significant**.

As a result, the output will not simply be a larger version of the input. The output will instead be **distorted**—a phenomenon known as **Intermodulation Distortion**.

**Q:** Good heavens! This sounds terrible. What exactly is **Intermodulation Distortion**, and what will it do to our signal output?!?

A: Say the input to the amplifier is sinusoidal, with magnitude a:

 $v_{in} = a \cos \omega t$ 

Jim Stiles

Using our knowledge of trigonometry, we can determine the result of the second term of the output Taylor series:

$$Bv_{in}^{2} = Ba^{2}\cos^{2}\omega t$$
$$= \frac{Ba^{2}}{2} + \frac{Ba^{2}}{2}\cos 2\omega t$$

We have created a harmonic of the input signal!

In other words, the input signal is at a frequency  $\omega$ , while the output includes a signal at **twice** that frequency ( $2\omega$ ).

We call this signal a **second** order product, as it is a result of **squaring** the input signal.

Note we also have a **cubed** term in the output signal equation:

$$\boldsymbol{v}_{out} = \boldsymbol{A}_{v} \boldsymbol{v}_{in} + \boldsymbol{B} \boldsymbol{v}_{in}^{2} + \boldsymbol{C} \boldsymbol{v}_{in}^{3} + \cdots$$

Using a trig identity, we find that:

$$C v_{in}^{3} = C a^{3} \cos^{3} \omega t$$
$$= \frac{C a^{3}}{2} \cos \omega t + \frac{C a^{3}}{4} \cos 3\omega t$$

Now we have produced a second harmonic (i.e.,  $3\omega$ )!

As you might expect, we call this harmonic signal a **third-order** product (since it's produced from  $v_{in}^3$ ).

**Q:** I confess that I am still a bit **befuddled**. You said that values B and C are typically **much** smaller that that of voltage gain  $A_v$ . Therefore it would seem that these harmonic signals would be **tiny** compared to the fundamental output signal  $A_v$  a cos  $\omega t$ . Thus, I **don't** why there's a problem!

To understand why intermodulation distortion can be a problem in amplifiers, we need to consider the **power** of the output signals.

We know that the power of a sinusoidal signal is proportional to its **magnitude squared**. Thus, we find that the power of each output signal is related to the input signal power as:

1rst-order output power  $\doteq P_1^{out} = A_v^2 P_{in} = G P_{in}$ 

2nd-order output power  $\doteq P_2^{out} = \frac{B^2}{4}P_{in}^2 = G_2P_{in}^2$ 

3rd-order output power  $\doteq P_3^{out} = \frac{C^2}{16}P_{in}^3 = G_3P_{in}^3$ 

where we have obviously defined  $G_2 \doteq B^2/4$  and  $G_3 \doteq C^2/16$ 

Note that unlike G, the values  $G_2$  and  $G_3$  are **not coefficients** (i.e., not unitless!). The value  $G_2$  obviously has units of inverse power (e.g.,  $mW^1$  or  $W^1$ ), while  $G_3$  has units of inverse power squared (e.g.,  $mW^2$  or  $W^2$ ).

We know that typically,  $G_2$  and  $G_3$  are much smaller than G. Thus, we are **tempted** to say that  $P_1^{out}$  is much larger than  $P_2^{out}$  or  $P_3^{out}$ .

But, we might be wrong !

**Q:** *Might* be wrong! Now I'm more confused than ever. Why can't we say **definitively** that the second and third order products are **insignificant**??

Look **closely** at the expressions for the output power of the first, second, and third order products:

 $P_1^{out} = G P_{in}$   $P_2^{out} = G_2 P_{in}^2$   $P_3^{out} = G_3 P_{in}^3$ 

This **first** order output power is of course **directly** proportional to the input power. However, the **second** order output power is proportional to the input power **squared**, while the **third** order output is proportional to the input power **cubed**!

Thus we find that if the input power is **small**, the second and third order products **are** insignificant. But, as the input power **increases**, the second and third order products get **big** in a hurry! For example, if we **double** the input power, the **first** order signal will of course likewise **double**. However, the **second** order power will **quadruple**, while the **third** order power will increase **8 times**.

For **large** input powers, the second and third order output products can in fact be **almost as large** as the first order signal!

Perhaps this can be most easily seen by expressing the above equations in **decibels**, e.g.,:

$$P_1^{out}(dBm) = G(dB) + P_{in}(dBm)$$

$$P_2^{out}(dBm) = G_2(dBm^{-1}) + 2[P_{in}(dBm)]$$

$$P_3^{out}(dBm) = G_3(dBm^{-2}) + 3[P_{in}(dBm)]$$

where we have used the fact that  $\log x^n = n \log x$ . Likewise, we have defined:

$$G_{2}(dBm^{-1}) = 10\log_{10}\left[\frac{G_{2}}{(1/1.0mW)}\right]$$
$$= 10\log_{10}[G_{2}(1.0mW)]$$

and:

### Hint: Just express everything in milliwatts!

Note the value  $2[P_{in}(dBm)]$  does **not** mean the value  $2P_{in}$  expressed in decibels. The value  $2[P_{in}(dBm)]$  is fact the value of  $P_{in}$  expressed in decibels—**times two**!

For **example**, if  $P_{in}(dBm) = -30 dBm$ , then  $2[P_{in}(dBm)] = -60 dBm$ . Likewise, if  $P_{in}(dBm) = 20 dBm$ , then  $2[P_{in}(dBm)] = 40 dBm$ .

What this means is that for every 1dB increase in input power  $P_{in}$  the fundamental (first-order) signal will increase 1dB; the second-order power will increase 2dB; and the third-order power will increase 3dB.

This is evident when we look at the three power equations (in decibels), as each is an equation of a line (i.e., y = m x + b).

For example, the equation:

$$P_3^{out}(dBm) = 3[P_{in}(dBm)] + G_3(dBm^{-2})$$
$$y = mx + b$$

describes a line with slope m = 3 and "y intercept"  $b = G_3(dBm^{-2})$  (where  $x = P_{in}(dBm)$  and  $y = P^{out}(dBm)$ ).



However, when the input power increases **beyond** 0 dBm (the right side of the plot), the second and third order products rapidly **catch up**! In fact, they will (theoretically) become **equal** to the first order product at some large input power.

The point at which each higher order product **equals** the firstorder signal is defined as the **intercept point**. Thus, we define the **second order intercept** point as the output power **when**:

$$P_2^{out} = P_1^{out} \doteq P_2^{int}$$
 Second - order intercept power

Likewise, the **third order intercept** point is defined as the third-order output power **when**:

$$P_3^{out} = P_1^{out} \doteq P_3^{int}$$
 Third - order intercept power

Using a little algebra you can show that:

$$P_2^{\text{int}} = \frac{G^2}{G_2}$$
 and  $P_3^{\text{int}} = \sqrt{\frac{G^3}{G_3}}$ 

Or, expressed in decibels:

$$P_2^{\text{int}}(dBm) = 2 \mathcal{G}(dB) - \mathcal{G}_2(dBm^{-1})$$

$$P_3^{\text{int}}(dBm) = \frac{3 G(dB) - G_3(dBm^{-2})}{2}$$

- \* Radio engineers specify the intermodulation distortion performance of a specific amplifier in terms of the **intercept points**, rather than values  $G_2$  and  $G_3$ .
- \* Generally, only the **third-order** intercept point is provided by amplifier manufactures (we'll see why later).
- Typical values of P<sub>3</sub><sup>int</sup> for a small-signal amplifier range
   from +20 dBm to +50 dBm
- \* Note that as  $G_2$  and  $G_3$  decrease, the intercept points increase.
- Therefore, the **higher** the intercept point of an amplifier, the better the amplifier !

One other **important point**: the **intercept points** for most amplifiers are much **larger** than the **compression point**! I.E.,:

 $P^{\text{int}} > P_{1dB}$ 

In other words the intercept points are "theoretical", in that we can **never**, in fact, increase the input power to the point that the higher order signals are **equal** to the fundamental signal power.

All signals, including the higher order signals, have a maximum limit that is determined by the amplifier **power supply**.

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### **Two-Tone Intermodulation**

**Q:** It doesn't seem to me that this **dad-gum** intermodulation distortion is really that much of a problem.

I mean, the first and second harmonics will likely be well outside the amplifier bandwidth, right?



A: True, the harmonics produced by intermodulation distortion typically are not a problem in radio system design. There is a problem, however, that is much worse than harmonic distortion!

This problem is called **two-tone** intermodulation distortion.

Say the input to an amplifier consists of **two** signals at **dissimilar** frequencies:

$$v_{in} = a \cos \omega_1 t + a \cos \omega_2 t$$

Here we will assume that both frequencies  $\omega_1$  and  $\omega_2$  are within the **bandwidth** of the amplifier, but are **not** equal to each other

 $(\omega_1 = \omega_2)$ .

This of course is a much more **realistic** case, as typically there will be **multiple** signals at the input to an amplifier!

For example, the two signals considered here could represent two **FM radio stations**, operating at frequencies within the FM band (i.e., 88.1 MHz  $\leq f_1 \leq 108.1$  MHz and 88.1 MHz  $\leq f_2 \leq 108.1$  MHz ).



**Q:** My point exactly! Intermodulation distortion will produce those **dog-gone** secondorder products:

$$\frac{a^2}{2}\cos 2\omega_1 t$$
 and  $\frac{a^2}{2}\cos 2\omega_2 t$ 

and gul-durn third order products:

$$\frac{a^3}{4}\cos 3\omega_1 t$$
 and  $\frac{a^3}{4}\cos 3\omega_2 t$ 

but these harmonic signals will lie well **outside** the FM band!

A: True! Again, the harmonic signals are not the problem. The problem occurs when the two input signals combine together to form additional second and third order products.

Recall an amplifier output is accurately described as:

$$\boldsymbol{v}_{out} = \boldsymbol{A}_{v_{in}} + \boldsymbol{B} \, \boldsymbol{v}_{in}^2 + \boldsymbol{C} \, \boldsymbol{v}_{in}^3 + \cdots$$

Consider first the **second-order** term if **two** signals are at the input to the amplifier:

$$v_2^{out} = B v_{in}^2$$
  
=  $B (a \cos \omega_1 t + a \cos \omega_2 t)^2$   
=  $B (a^2 \cos^2 \omega_1 t + 2a^2 \cos \omega_1 t \cos \omega_2 t + a^2 \cos^2 \omega_2 t)$ 

Note the first and third terms of the above expression are **precisely** the same as the terms we examined on the previous handout. They result in **harmonic** signals at frequencies  $2\omega_1$  and  $2\omega_2$ , respectively.

The **middle** term, however, is something **new**. Note **it** involves the product of  $\cos \omega_1 t$  and  $\cos \omega_2 t$ . Again using our knowledge of **trigonometry**, we find:

$$2a^2\cos\omega_1 t \ \cos\omega_2 t = a^2\cos(\omega_2 - \omega_1)t + a^2\cos(\omega_2 + \omega_1)t$$

Note that since  $\cos(-x) = \cos x$ , we can **equivalently** write this as:

$$2a^2\cos\omega_1 t \ \cos\omega_2 t = a^2\cos(\omega_1 - \omega_2)t + a^2\cos(\omega_1 + \omega_2)t$$

Either way, the result is obvious—we produce two new signals!

These new second-order signals oscillate at frequencies  $(\omega_1 + \omega_2)$  and  $|\omega_1 - \omega_2|$ .

Thus, if we looked at the **frequency spectrum** (i.e., signal power as a function of frequency) of an amplifier **output** when two sinusoids are at the input, we would see something like this:



Note that the new terms have a frequency that is either much higher than both  $\omega_1$  and  $\omega_2$  (i.e.,  $(\omega_1 + \omega_2)$ ), or much lower than both  $\omega_1$  and  $\omega_2$  (i.e.,  $|\omega_1 - \omega_2|$ ).

Either way, these new signals will typically be **outside** the amplifier bandwidth!

**Q:** I thought you said these "two-tone" intermodulation products were some "**big problem**". These sons of a gun appear to be **no more** a problem than the harmonic signals!



A: This observation is indeed correct for **second**-order, twotone intermodulation products. But, we have **yet** to examine the **third**-order terms! I.E.,

$$v_3^{out} = C v_{in}^3$$
$$= C (a \cos \omega_1 t + a \cos \omega_2 t)^3$$

If we multiply this all out, and again apply our trig knowledge, we find that a **bunch** of new **third-order** signals are created.

Among these signals, of course, are the second harmonics  $\cos 3\omega_1 t$  and  $\cos 3\omega_2 t$ . Additionally, however, we get these new signals:

$$\cos(2\omega_2 - \omega_1)t$$
 and  $\cos(2\omega_1 - \omega_2)t$ 

Note since  $\cos(-x) = \cos x$ , we can **equivalently** write these terms as:

 $\cos(\omega_1 - 2\omega_2)t$  and  $\cos(\omega_2 - 2\omega_1)t$ 

Either way, it is apparent that the **third-order** products include signals at frequencies  $|\omega_1 - 2\omega_2|$  and  $|\omega_2 - 2\omega_1|$ .

Now lets look at the output spectrum with **these new** thirdorder products included:



Now you should see the problem! These third-order products are very close in frequency to  $\omega_1$  and  $\omega_2$ . They will likely lie within the bandwidth of the amplifier!

For example, if  $f_1$ =100 MHz and  $f_2$ =101 MHz, then  $2f_2 - f_1$ =102 MHz and  $2f_1 - f_2$ = 99 MHz. All frequencies are **well** within the FM radio bandwidth!

Thus, these **third-order**, **two-tone** intermodulation products are the **most significant** distortion terms.

This is why we are most concerned with the **third-order** intercept point of an amplifier!



I only select amplifiers with the **highest possible** 3<sup>rd</sup>order intercept point!

# The Amplifier Spec Sheet

Here's a list of some of the most **important** amplifier specifications:

### <u>Gain (dB)</u>

Provides the "nominal" ( $\leftarrow$  a vendor weasel word) gain within the bandwidth of the amplifier.

### Gain Flatness(dB)

Usually expressed as +/- (e.g. +/- 0.5 dB), the specification indicates how much the gain varies across the "middle" portion of the amplifier bandwidth.

### Bandwidth (Hz)

Expressed in terms of a low-end frequency (i.e.,  $f_L$ ) and a highend frequency (i.e.,  $f_H$ ), this value is generally specified as the frequencies where the gain has reduced to 3dB less than its nominal value.

### **Input Impedance**

Can be expressed in a number of ways:  $S_{11}$ ,  $Z_{in}$ ,  $\Gamma$ , return loss, or VSWR.

### <u>Ouput Impedance</u>

Can be expressed in a number of ways: S\_{22}, Z\_{in},  $\Gamma$ , return loss, or VSWR.

### **Reverse Isolation (dB)**

The larger the number, the better.

### D.C. Power

Simply the product of the DC current and DC voltage.

### <u>1 dB compression point (Watts, dBm, dBw)</u>

Regarded as the maximum **output power** the amplifier can produce.

### <u>3<sup>rd</sup> order intercept point (Watts, dBm, dBw)</u>

The larger the number, the better.

### <u>Noise Figure (dB)</u>

We will learn about this later!