<u>A Transmission Line</u> <u>Connecting Source & Load</u>

We can think of a transmission line as a conduit that allows **power** to flow **from** an **output** of one device/network **to** an **input** of another.

To simplify our analysis, we can model the **input** of the device **receiving** the power with it input impedance (e.g., Z_L), while we can model the device **output delivering** the power with its Thevenin's or Norton's equivalent circuit.



 V_{g}

Typically, the power source is modeled with its **Thevenin's** equivalent; however, we will find that the **Norton's** equivalent circuit is useful if we express the remainder of the circuit in terms of its **admittance** values (e.g., $Y_0, Y_L, Y(z)$).

 Z_0

 I_i

+

 V_i

 $\mathbf{Z} = -\ell$

 Z_{g}

Recall from the telegrapher's equations that the current and voltage along the transmission line are:

$$V(z) = V_0^+ e^{-j\beta z} + V_0^- e^{+j\beta z}$$

$$I(z) = \frac{V_0^+}{Z_0} e^{-j\beta z} - \frac{V_0^-}{Z_0} e^{+j\beta z}$$

At z = 0, we enforced the **boundary condition** resulting from Ohm's Law:

$$Z_{L} = \frac{V_{L}}{I_{L}} = \frac{V(z=0)}{I(z=0)} = \frac{\left(V_{0}^{+} + V_{0}^{-}\right)}{\left(\frac{V_{0}^{+}}{Z_{0}} - \frac{V_{0}^{-}}{Z_{0}}\right)}$$

 Z_{L}

z = 0



$$V(z) = V_0^+ \left[e^{-j\beta z} + \Gamma_L e^{+j\beta z} \right]$$

$$I(z) = \frac{V_0^+}{Z_0} \left[e^{-j\beta z} - \Gamma_L e^{+j\beta z} \right]$$

We are left with the **question**: just what is the **value** of complex constant V_0^+ ?!?

This constant depends on the signal source! To determine its exact value, we must now apply boundary conditions at $z = -\ell$.

We know that at the **beginning** of the transmission line:

$$V(z = -\ell) = V_0^+ \left[e^{+j\beta\ell} + \Gamma_L e^{-j\beta\ell} \right]$$

$$I(z = -\ell) = \frac{V_0^+}{Z_0} \left[e^{+j\beta\ell} - \Gamma_L e^{-j\beta\ell} \right]$$

Likewise, we know that the source must satisfy:

$$V_g = V_i + Z_g I_i$$

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$$V_{0}^{+} = V_{g} e^{-j\beta\ell} \frac{Z_{0}}{Z_{0} (1 + \Gamma_{in}) + Z_{g} (1 - \Gamma_{in})}$$

where:

$$\Gamma_{in} = \Gamma(\boldsymbol{z} = -\ell) = \Gamma_{I} \boldsymbol{e}^{-j^{2}\beta\ell}$$

Note this result looks different than the equation in your textbook (eq. 2.71):

$$V_0^+ = V_g \frac{Z_0}{Z_0 + Z_g} \frac{e^{-j\beta\ell}}{\left(1 - \Gamma_L \Gamma_g e^{-j2\beta\ell}\right)}$$

where:

$$\Gamma_g \doteq \frac{Z_g - Z_0}{Z_g + Z_0}$$

I like my expression better.

Although the two equations are equivalent, **my** expression is explicitly written in terms of $\Gamma_{in} = \Gamma(z = -\ell)$ (a very **useful**, **precise**, and **unambiguous** value), while the book's expression is written in terms of this **so-called** "source reflection coefficient" Γ_g (a **misleading**, **confusing**, **ambiguous**, and mostly **useless** value).

Specifically, we might be **tempted** to equate Γ_g with the value $\Gamma(z = -\ell) = \Gamma_m$, but it is **not** $(\Gamma_g \neq \Gamma(z = -\ell))!$

There is one **very important** point that must be made about the result:

$$V_{0}^{+} = V_{g} e^{-j\beta\ell} \frac{Z_{0}}{Z_{0} \left(1 + \Gamma_{in}\right) + Z_{g} \left(1 - \Gamma_{in}\right)}$$

And that is—the wave $V_0^+(z)$ incident on the load Z_L is actually dependent on the value of load Z_L !!!!!

Remember:

$$\Gamma_{in} = \Gamma(\boldsymbol{z} = -\ell) = \Gamma_L \, \boldsymbol{e}^{-j2\beta\ell}$$

We tend to think of the incident wave $V_0^+(z)$ being "caused" by the source, and it is certainly true that $V_0^+(z)$ depends on the source—after all, $V_0^+(z) = 0$ if $V_g = 0$. However, we find from the equation above that it **likewise** depends on the value of the load!

Thus we **cannot**—in general—consider the incident wave to be the "**cause**" and the reflected wave the "**effect**". Instead, each wave must obtain the proper **amplitude** (e.g., V_0^+, V_0^-) so that the boundary conditions are satisfied at **both** the beginning and end of the transmission line.