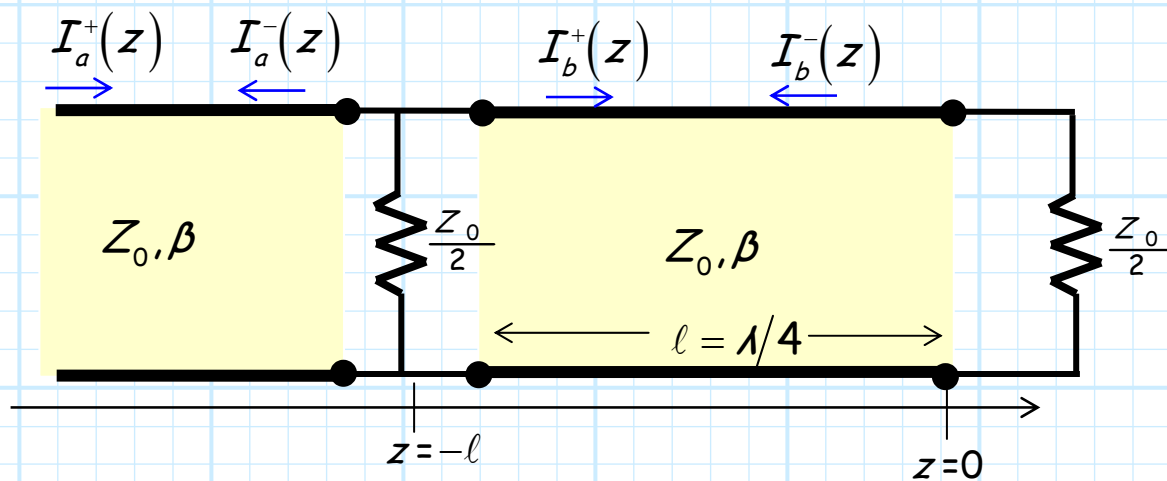


Example: Another Boundary Condition Problem



The **total** voltage along the transmission line shown above is expressed as:

$$V(z) = \begin{cases} V_{0a}^+ e^{-j\beta z} + V_{0a}^- e^{+j\beta z} & z < -l \\ V_{0b}^+ e^{-j\beta z} + V_{0b}^- e^{+j\beta z} & -l < z < 0 \end{cases}$$

Carefully determine and apply boundary conditions at both $z = 0$ and $z = -l$ to find the three values:

$$\frac{V_{0a}^-}{V_{0a}^+}, \quad \frac{V_{0b}^+}{V_{0a}^+}, \quad \frac{V_{0b}^-}{V_{0a}^+}$$

Solution

From the telegrapher's equation, we likewise know that the current along the transmission lines is:

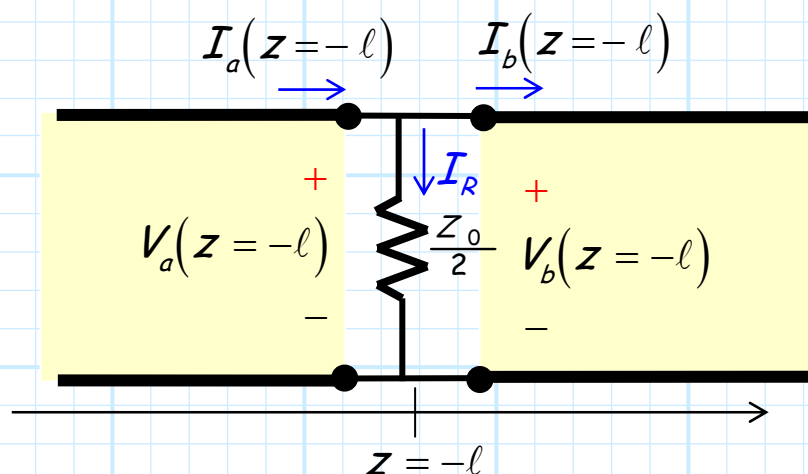
$$I(z) = \begin{cases} \frac{V_{0a}^+}{Z_0} e^{-j\beta z} - \frac{V_{0a}^-}{Z_0} e^{+j\beta z} & z < -l \\ \frac{V_{0b}^+}{Z_0} e^{-j\beta z} - \frac{V_{0b}^-}{Z_0} e^{+j\beta z} & -l < z < 0 \end{cases}$$

To find the values:

$$\frac{V_{0a}^-}{V_{0a}^+}, \quad \frac{V_{0b}^+}{V_{0a}^+}, \quad \frac{V_{0b}^-}{V_{0a}^+}$$

We need only to evaluate boundary conditions!

Boundary Conditions at $z = -l$



From KVL, we conclude:

$$V_a(z = -l) = V_b(z = -l)$$

From KCL:

$$I_a(z = -\ell) = I_b(z = -\ell) + I_R$$

And from Ohm's Law:

$$I_R = \frac{V_a(z = -\ell)}{Z_0/2} = \frac{2V_a(z = -\ell)}{Z_0} = \frac{2V_b(z = -\ell)}{Z_0}$$

We likewise know from the telegrapher's equation that:

$$\begin{aligned} V_a(z = -\ell) &= V_{0a}^+ e^{-j\beta(-\ell)} + V_{0a}^- e^{+j\beta(-\ell)} \\ &= V_{0a}^+ e^{+j\beta\ell} + V_{0a}^- e^{-j\beta\ell} \end{aligned}$$

And since $\ell = \lambda/4$, we find:

$$\beta\ell = \left(\frac{2\pi}{\lambda}\right)\left(\frac{\lambda}{4}\right) = \frac{\pi}{2}$$

And so:

$$\begin{aligned} V_a(z = -\ell) &= V_{0a}^+ e^{+j\beta\ell} + V_{0a}^- e^{-j\beta\ell} \\ &= V_{0a}^+ e^{+j(\pi/2)} + V_{0a}^- e^{-j(\pi/2)} \\ &= V_{0a}^+ (j) + V_{0a}^- (-j) \\ &= j(V_{0a}^+ - V_{0a}^-) \end{aligned}$$

We similarly find that:

$$V_b(z = -\ell) = j(V_{0b}^+ - V_{0b}^-)$$

and for currents:

$$I_a(z = -\ell) = j \frac{V_{0a}^+ + V_{0a}^-}{Z_0}$$

$$I_b(z = -\ell) = j \frac{V_{0b}^+ + V_{0b}^-}{Z_0}$$

Inserting these results into our KVL boundary condition statement:

$$\begin{aligned} V_a(z = -\ell) &= V_b(z = -\ell) \\ j(V_{0a}^+ - V_{0a}^-) &= j(V_{0b}^+ - V_{0b}^-) \\ V_{0a}^+ - V_{0a}^- &= V_{0b}^+ - V_{0b}^- \end{aligned}$$

Normalizing to (i.e., dividing by) V_{0a}^+ , we conclude:

$$1 - \frac{V_{0a}^-}{V_{0a}^+} = \frac{V_{0b}^+}{V_{0a}^+} - \frac{V_{0b}^-}{V_{0a}^+}$$

From Ohm's Law:

$$I_R = \frac{2V_a(z = -\ell)}{Z_0} = \frac{2j(V_{0a}^+ - V_{0a}^-)}{Z_0}$$

$$I_R = \frac{2V_b(z = -\ell)}{Z_0} = \frac{2j(V_{0b}^+ - V_{0b}^-)}{Z_0}$$

And finally from our KCL boundary condition:

$$\begin{aligned} I_a(z = -\ell) &= I_b(z = -\ell) + I_R \\ j \frac{V_{0a}^+ + V_{0a}^-}{Z_0} &= j \frac{V_{0b}^+ + V_{0b}^-}{Z_0} + I_R \end{aligned}$$

After an **enjoyable** little bit of algebra, we can thus conclude:

$$V_{0a}^+ + V_{0a}^- = V_{0b}^+ + V_{0b}^- - j I_R Z_0$$

And inserting the result from Ohm's Law:

$$\begin{aligned} V_{0a}^+ + V_{0a}^- &= V_{0b}^+ + V_{0b}^- - j I_R Z_0 \\ &= V_{0b}^+ + V_{0b}^- - j \left(\frac{2j(V_{0b}^+ - V_{0b}^-)}{Z_0} \right) Z_0 \\ &= V_{0b}^+ + V_{0b}^- - 2j^2 (V_{0b}^+ - V_{0b}^-) \left(\frac{Z_0}{Z_0} \right) \\ &= V_{0b}^+ + V_{0b}^- - 2(-1)(V_{0b}^+ - V_{0b}^-) \\ &= V_{0b}^+ + V_{0b}^- + 2V_{0b}^+ - 2V_{0b}^- \\ &= 3V_{0b}^+ - V_{0b}^- \end{aligned}$$

Again normalizing to V_{0a}^+ , we get a second relationship:

$$1 + \frac{V_{0a}^-}{V_{0a}^+} = 3 \frac{V_{0b}^+}{V_{0a}^+} - \frac{V_{0b}^-}{V_{0a}^+}$$

Q: But wait! We now have **two** equations:

$$1 - \frac{V_{0a}^-}{V_{0a}^+} = \frac{V_{0b}^+}{V_{0a}^+} - \frac{V_{0b}^-}{V_{0a}^+} \qquad 1 + \frac{V_{0a}^-}{V_{0a}^+} = 3 \frac{V_{0b}^+}{V_{0a}^+} - \frac{V_{0b}^-}{V_{0a}^+}$$

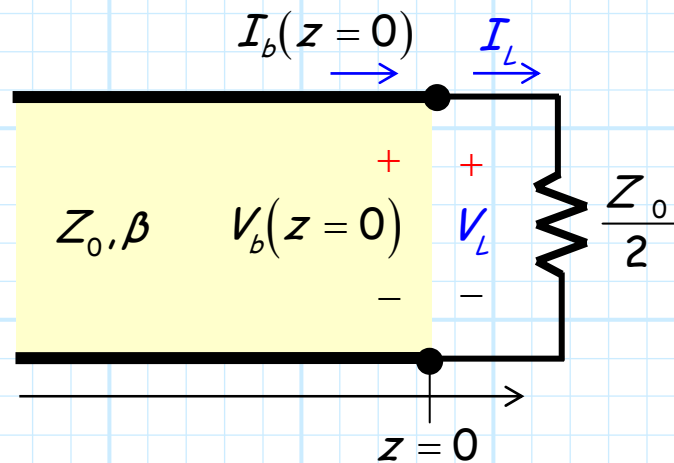
but **three** unknowns:

$$\frac{V_{0a}^-}{V_{0a}^+}, \frac{V_{0b}^+}{V_{0a}^+}, \frac{V_{0b}^-}{V_{0a}^+}$$

Did we make a *mistake* somewhere?

A: Nope! We just have more work to do. After all, there is a yet **another** boundary to be analyzed!

Boundary Conditions at $z = 0$



From KVL, we conclude:

$$V_b(z=0) = V_L$$

From KCL:

$$I_b(z=0) = I_L$$

And from Ohm's Law:

$$I_L = \frac{V_L}{Z_0/2} = \frac{2V_L}{Z_0}$$

We likewise know from the telegrapher's equation that:

$$\begin{aligned} V_b(z=0) &= V_{0b}^+ e^{-j\beta(0)} + V_{0b}^- e^{+j\beta(0)} \\ &= V_{0b}^+ (1) + V_{0b}^- (1) \\ &= V_{0b}^+ + V_{0b}^- \end{aligned}$$

We similarly find that:

$$I_b(z=0) = \frac{V_{0b}^+ - V_{0b}^-}{Z_0}$$

Combing this with the above results:

$$I_L = \frac{2V_L}{Z_0}$$

$$I_b(z=0) = \frac{2V_b(z=0)}{Z_0}$$

$$\frac{V_{0b}^+ - V_{0b}^-}{Z_0} = \frac{2(V_{0b}^+ + V_{0b}^-)}{Z_0}$$

From which we conclude:

$$V_{0b}^+ - V_{0b}^- = 2(V_{0b}^+ + V_{0b}^-) \Rightarrow -3V_{0b}^- = V_{0b}^+$$

And so:

$$V_{0b}^- = -\frac{1}{3}V_{0b}^+$$

Note that we could have also determined this using the load reflection coefficient:

$$\frac{V_b^-(z=0)}{V_b^+(z=0)} = \Gamma(z=0) = \Gamma_0$$

Where:

$$V_b^-(z=0) = V_{0b}^- e^{+j\beta(0)} = V_{0b}^-$$

$$V_b^+(z=0) = V_{0b}^+ e^{-j\beta(0)} = V_{0b}^+$$

And we use the boundary condition:

$$\Gamma_{0b} = \Gamma_{Lb} = \frac{Z_L - Z_0}{Z_L + Z_0} = \frac{0.5Z_0 - Z_0}{0.5Z_0 + Z_0} = \frac{-0.5}{1.5} = -\frac{1}{3}$$

Therefore, we arrive at the **same result** as before:

$$\frac{V_b^-(z=0)}{V_b^+(z=0)} = \Gamma_{0b}$$

$$\frac{V_{0b}^-}{V_{0b}^+} = -\frac{1}{3}$$

Either way, we can use this result to simplify our first set of boundary conditions:

$$\begin{aligned} 1 - \frac{V_{0a}^-}{V_{0a}^+} &= \frac{V_{0b}^+}{V_{0a}^+} - \frac{V_{0b}^-}{V_{0a}^+} \\ &= \frac{V_{0b}^+}{V_{0a}^+} - \frac{-V_{0b}^+/3}{V_{0a}^+} \\ &= \frac{V_{0b}^+}{V_{0a}^+} + \frac{1}{3} \frac{V_{0b}^+}{V_{0a}^+} \\ &= \frac{4}{3} \frac{V_{0b}^+}{V_{0a}^+} \end{aligned}$$

And:

$$\begin{aligned}
 1 + \frac{V_{0a}^-}{V_{0a}^+} &= 3 \frac{V_{0b}^+}{V_{0a}^+} - \frac{V_{0b}^-}{V_{0a}^+} \\
 &= 3 \frac{V_{0b}^+}{V_{0a}^+} - \frac{-V_{0b}^+}{3} \\
 &= 3 \frac{V_{0b}^+}{V_{0a}^+} + \frac{1}{3} \frac{V_{0b}^+}{V_{0a}^+} \\
 &= \frac{10}{3} \frac{V_{0b}^+}{V_{0a}^+}
 \end{aligned}$$

NOW we have **two** equations and **two** unknowns:

$$1 - \frac{V_{0a}^-}{V_{0a}^+} = \frac{4}{3} \frac{V_{0b}^+}{V_{0a}^+} \qquad 1 + \frac{V_{0a}^-}{V_{0a}^+} = \frac{10}{3} \frac{V_{0b}^+}{V_{0a}^+}$$

Adding the two equations, we find:

$$\begin{aligned}
 \left(1 - \frac{V_{0a}^-}{V_{0a}^+} \right) + \left(1 + \frac{V_{0a}^-}{V_{0a}^+} \right) &= \left(\frac{4}{3} \frac{V_{0b}^+}{V_{0a}^+} \right) + \left(\frac{10}{3} \frac{V_{0b}^+}{V_{0a}^+} \right) \\
 2 &= \frac{14}{3} \frac{V_{0b}^+}{V_{0a}^+} \\
 \frac{3}{7} &= \frac{V_{0b}^+}{V_{0a}^+}
 \end{aligned}$$

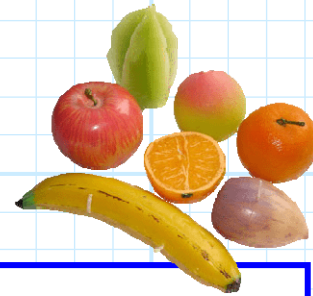
And so using the second equation above:

$$\begin{aligned}\frac{V_{0a}^-}{V_{0a}^+} &= \frac{10}{3} \frac{V_{0b}^+}{V_{0a}^+} - 1 \\ &= \frac{10}{3} \frac{3}{7} - 1 \\ &= \frac{3}{7}\end{aligned}$$

And finally, from one of our original boundary conditions:

$$\begin{aligned}\frac{V_{0b}^-}{V_{0a}^+} &= \frac{V_{0b}^+}{V_{0a}^+} - 1 + \frac{V_{0a}^-}{V_{0a}^+} \\ &= \frac{3}{7} - 1 + \frac{3}{7} \\ &= -\frac{1}{7}\end{aligned}$$

And so now we **summarize** the fruit of our labor:



$$\frac{V_{0a}^-}{V_{0a}^+} = \frac{3}{7} \quad \frac{V_{0b}^+}{V_{0a}^+} = \frac{3}{7} \quad \frac{V_{0b}^-}{V_{0a}^+} = -\frac{1}{7}$$

Yes it is! It's time for a **sanity check!!!**

The first of our boundary condition equations:

$$1 - \frac{V_{0a}^-}{V_{0a}^+} = \frac{V_{0b}^+}{V_{0a}^+} - \frac{V_{0b}^-}{V_{0a}^+}$$

$$1 - \frac{3}{7} = \frac{3}{7} - \left(-\frac{1}{7} \right)$$

$$\frac{4}{7} = \frac{4}{7} \quad \checkmark$$

And from the second:

$$1 + \frac{V_a^-}{V_a^+} = 3 \frac{V_b^+}{V_a^+} - \frac{V_b^-}{V_a^+}$$

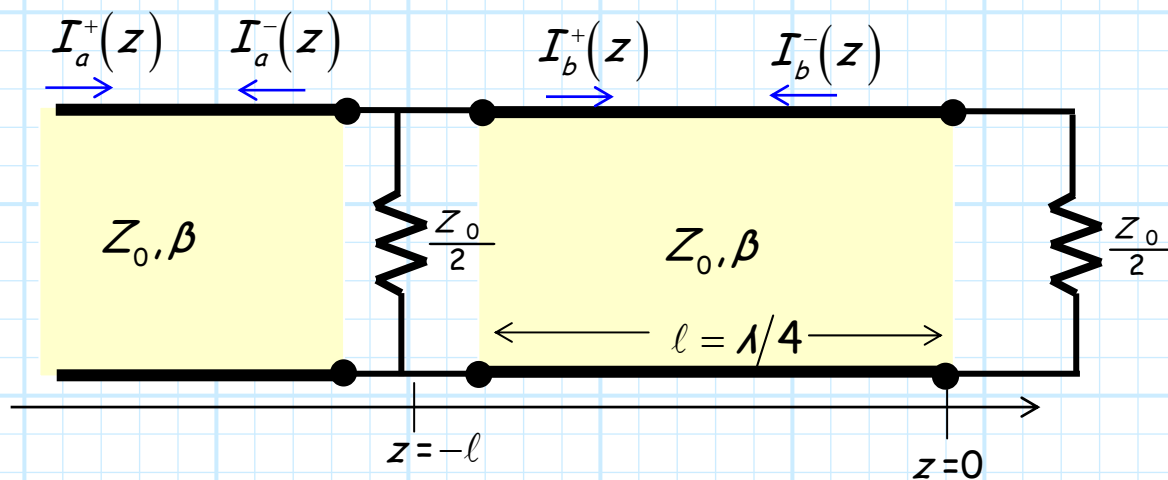
$$1 + \frac{3}{7} = 3 \frac{3}{7} - \left(-\frac{1}{7} \right)$$

$$\frac{10}{7} = \frac{10}{7} \quad \checkmark$$

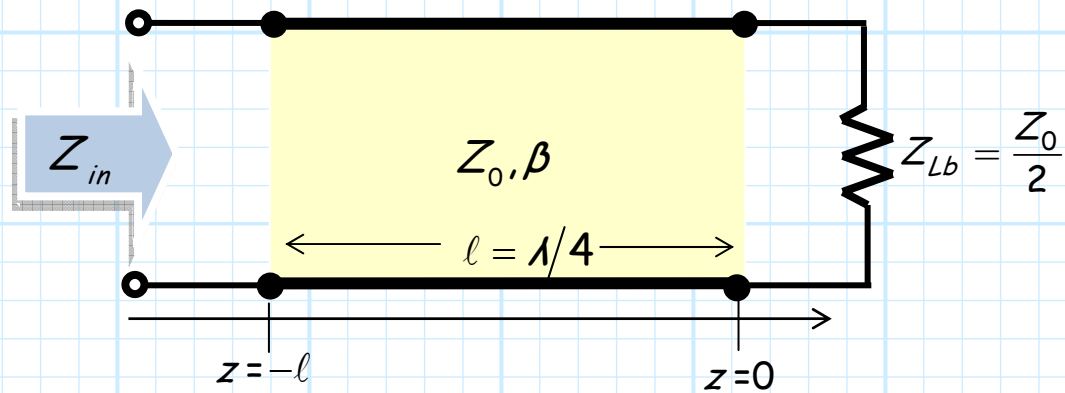
Notice that we can also verify the result:

$$\frac{V_{0a}^-}{V_{0a}^+} = \frac{3}{7}$$

By using the equivalent circuit of:



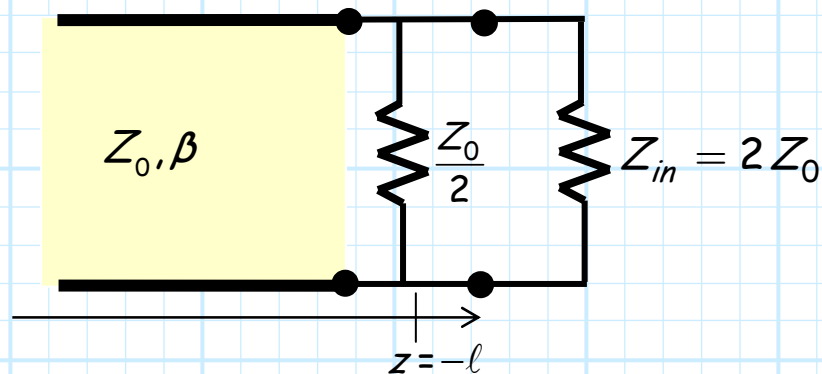
Specifically, we can determine the input impedance of this circuit:



Since the transmission line is the **special case** of one quarter wavelength, we know that:

$$Z_{in} = \frac{Z_0^2}{0.5Z_0} = 2.0 Z_0$$

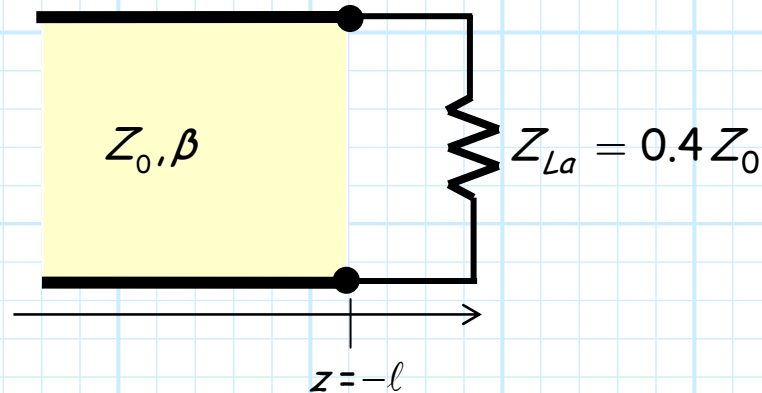
And so the equivalent circuit becomes:



Where the two parallel impedances combine as:

$$0.5Z_0 \parallel 2Z_0 = \frac{Z_0}{2.5} = 0.4Z_0$$

And so the equivalent load at $z = -l$ is $0.4Z_0$:



Now, the reflection coefficient of **this** load is:

$$\Gamma_{La} = \frac{0.4Z_0 - Z_0}{0.4Z_0 + Z_0} = \frac{-0.6}{1.4} = -\frac{3}{7}$$

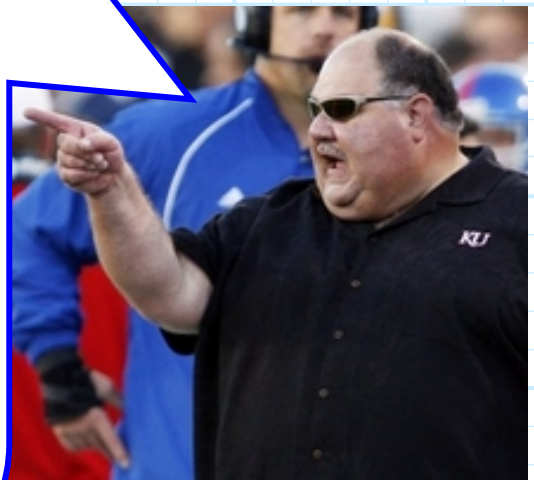
Q: *Wait a second! Using your fancy "boundary conditions" to solve the problem, you **earlier** arrived at the conclusion:*

$$\frac{V_{0a}^-}{V_{0a}^+} = \frac{3}{7}$$

But now we find that instead:

$$\frac{V_{0a}^-}{V_{0a}^+} = \Gamma_{La} = -\frac{3}{7}$$

Apparently your annoyingly pretentious boundary condition analysis introduced some sort of sign error !



A: Absolutely not! The boundary condition analysis is perfectly correct, and:

$$\frac{V_{0a}^-}{V_{0a}^+} = \frac{3}{7}$$

is the right answer.

The statement:

~~$$\frac{V_a^-}{V_a^+} = \Gamma_{La} = -\frac{3}{7}$$~~



is **erroneous!**

Q: *But how could that possibly be? You earlier determined that:*

$$\frac{V_{0b}^-}{V_{0b}^+} = \Gamma_{Lb} = -\frac{1}{3}$$

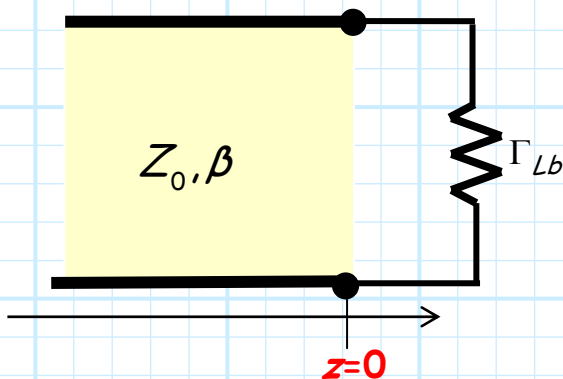
So why then is:

$$\frac{V_{0a}^-}{V_{0a}^+} \neq \Gamma_{La} \quad \text{????}$$



A: In the first case, load Γ_{Lb} is located at position $z = 0$, so that:

$$\Gamma_{Lb} = \frac{V_b^-(z=0)}{V_b^+(z=0)} = \frac{V_b^-}{V_b^+}$$

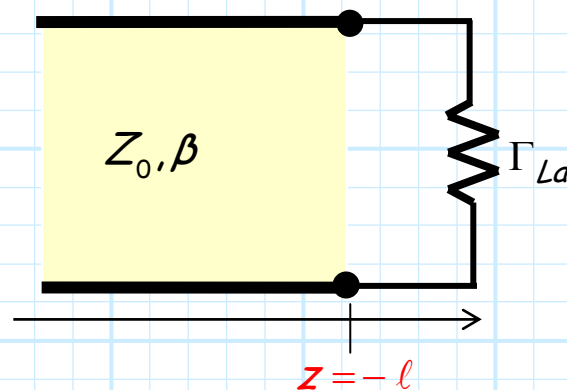


Note this result can be more compactly stated as a boundary condition requirement:

$$\Gamma_{Lb} = \Gamma(z = 0) = \frac{V_{0b}^-}{V_{0b}^+} e^{+j\beta(0)} = \frac{V_{0b}^-}{V_{0b}^+} = \Gamma_{0b}$$

For the **second** case, the load Γ_{Lb} is located **instead** at position $z = -\ell$, so that:

$$\Gamma_{La} = \frac{V_a^-(z = -\ell)}{V_a^+(z = -\ell)} = \frac{V_{0a}^- e^{-j\beta\ell}}{V_{0a}^+ e^{+j\beta\ell}} = \frac{V_{0a}^-}{V_{0a}^+} e^{-j2\beta\ell} = \Gamma_{0a} e^{-j2\beta\ell}$$



Note this result can be more compactly stated as a boundary condition requirement:

$$\Gamma_{La} = \Gamma(z = -\ell) = \frac{V_{0a}^-}{V_{0a}^+} e^{-j2\beta\ell}$$

From the equation above we find:



$$\frac{V_{0a}^-}{V_{0a}^+} = \Gamma_{La} e^{+j2\beta\ell} = -\frac{3}{7} e^{+j\pi} = +\frac{3}{7}$$

*That's **precisely** the same result as we determined earlier using our **boundary conditions!***

*Our answers are **good!***