

Special Cases of Source and Load Impedance

Let's look at **specific cases** of Z_g and Z_L , and determine how they affect V_0^+ and P_{abs} .

$$Z_g = Z_0$$

For this case, we find that V_0^+ **simplifies** greatly:

$$\begin{aligned} V_0^+ &= V_g e^{-j\beta\ell} \frac{Z_0}{Z_0(1 + \Gamma_{in}) + Z_g(1 - \Gamma_{in})} \\ &= V_g e^{-j\beta\ell} \frac{Z_0}{Z_0(1 + \Gamma_{in}) + Z_0(1 - \Gamma_{in})} \\ &= V_g e^{-j\beta\ell} \frac{1}{1 + \Gamma_{in} + 1 - \Gamma_{in}} \\ &= \frac{1}{2} V_g e^{-j\beta\ell} \end{aligned}$$

Look at what **this** says!

It says that the incident wave in this case is **independent** of the load attached at the other end!

Thus, for the **one** case $Z_g = Z_0$, we in fact can consider $V^+(z)$ as being the source wave, and then the reflected wave $V^-(z)$ as being the result of this stimulus.

Remember, the complex value V_0^+ is the value of the incident wave evaluated at the end of the transmission line ($V_0^+ = V^+(z=0)$). We can likewise determine the value of the incident wave at the **beginning** of the transmission line (i.e., $V^+(z=-\ell)$). For this case, where $Z_g = Z_0$, we find that this value can be very simply stated (!):

$$\begin{aligned} V^+(z=-\ell) &= V_0^+ e^{-j\beta(z=-\ell)} \\ &= \left(\frac{1}{2} V_g e^{-j\beta\ell} \right) e^{+j\beta\ell} \\ &= \frac{V_g}{2} \end{aligned}$$

Likewise, we find that the delivered power for this case can be simply stated as:

$$\begin{aligned} P_{abs} &= \frac{|V_0^+|^2}{2 Z_0} (1 - |\Gamma_L|^2) \\ &= \frac{|V_g|^2}{8 Z_0} (1 - |\Gamma_L|^2) \end{aligned}$$

$$Z_L = Z_0$$

In this case, we find that $\Gamma_L = 0$, and thus $\Gamma_{in} = 0$. As a result:

$$\begin{aligned} V_0^+ &= V_g e^{-j\beta\ell} \frac{Z_0}{Z_0(1 + \Gamma_{in}) + Z_g(1 - \Gamma_{in})} \\ &= V_g e^{-j\beta\ell} \frac{Z_0}{Z_0 + Z_g} \end{aligned}$$

Likewise, we find that:

$$P_{abs} = \frac{|V_0^+|^2}{2Z_0} (1 - |\Gamma_L|^2) = \frac{|V_0^+|^2}{2Z_0}$$

Here the delivered power P_{abs} is simply that of the incident wave (P^+), as the matched condition causes the reflected power to be zero ($P^- = 0$)!

Inserting the value of V_0^+ , we find:

$$\begin{aligned}
 P_{abs} &= \frac{|V_0^+|^2}{2 Z_0} \\
 &= \frac{|V_g|^2}{2 Z_0} \frac{(Z_0)^2}{|Z_0 + Z_g|^2} \\
 &= \frac{|V_g|^2}{2} \frac{Z_0}{|Z_0 + Z_g|^2}
 \end{aligned}$$

Note that this result can likewise be found by recognizing that $Z_{in} = Z_0$ when $Z_L = Z_0$:

$$\begin{aligned}
 P_{abs} &= \frac{1}{2} \frac{|V_g|^2}{|Z_g + Z_{in}|^2} \operatorname{Re}\{Z_{in}\} \\
 &= \frac{1}{2} \frac{|V_g|^2}{|Z_g + Z_0|^2} Z_0 \\
 &= \frac{|V_g|^2}{2} \frac{Z_0}{|Z_g + Z_0|^2}
 \end{aligned}$$

$$Z_{in} = Z_g^*$$

For this case, we find Z_L takes on whatever value required to make $Z_{in} = Z_g^*$. This is a **very** important case!

First, using the fact that:

$$\Gamma_{in} = \frac{Z_{in} - Z_0}{Z_{in} + Z_0} = \frac{Z_g^* - Z_0}{Z_g^* + Z_0}$$

We can show that (trust me!):

$$V_0^+ = V_g e^{-j\beta\ell} \frac{Z_g^* + Z_0}{4\text{Re}\{Z_g\}}$$

Not a particularly interesting result, but let's look at the absorbed power.

$$\begin{aligned} P_{abs} &= \frac{1}{2} \frac{|V_g|^2}{|Z_g + Z_{in}|^2} \text{Re}\{Z_{in}\} \\ &= \frac{1}{2} \frac{|V_g|^2}{|Z_g + Z_g^*|^2} \text{Re}\{Z_g^*\} \\ &= \frac{1}{2} \frac{|V_g|^2}{|2\text{Re}\{Z_g^*\}|^2} \text{Re}\{Z_g^*\} \\ &= \frac{1}{2} |V_g|^2 \frac{1}{4\text{Re}\{Z_g^*\}} \doteq P_{avl} \end{aligned}$$

Although this result does not look particularly interesting **either**, we find the result is **very** important!

It can be shown that—for a **given** V_g and Z_g —the value of input impedance Z_{in} that will absorb the **largest possible** amount of power is the value $Z_{in} = Z_g^*$.

This case is known as the **conjugate match**, and is essentially the goal of every transmission line problem—to deliver the largest possible power to Z_{in} , and thus to Z_L as well!

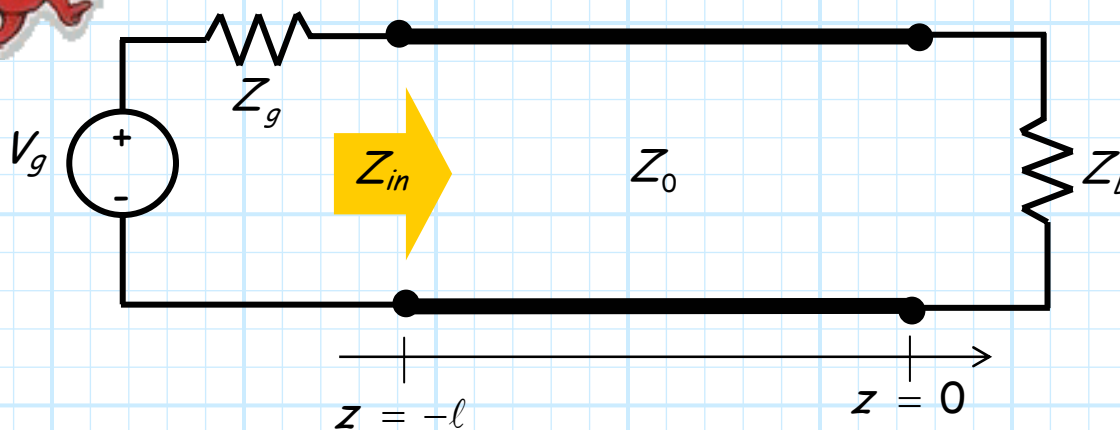
This maximum delivered power is known as the **available power** (P_{avl}) of the source.

There are **two** very important things to understand about this result!



Very Important Thing #1

Consider again the terminated transmission line:



Recall that if $Z_L = Z_0$, the **reflected** wave will be **zero**, and the absorbed power will be:

$$P_{abs} = \frac{|V_g|^2}{2} \frac{Z_0}{|Z_0 + Z_g|^2} \leq P_{avl}$$

But note if $Z_L = Z_0$, the input impedance $Z_{in} = Z_0$ —but then $Z_{in} \neq Z_g^*$ (generally)! In other words, $Z_L = Z_0$ does **not** (generally) result in a **conjugate match**, and thus setting $Z_L = Z_0$ does **not** result in maximum power absorption!

Q: *Huh!? This makes **no** sense! A load value of $Z_L = Z_0$ will **minimize** the reflected wave ($P^- = 0$)—**all** of the incident power will be absorbed.*

*Any other value of $Z_L = Z_0$ will result in **some** of the incident wave being reflected—how in the world could this **increase** absorbed power?*

*After all, just **look** at the expression for absorbed power:*

$$P_{abs} = \frac{|V_0^+|^2}{2 Z_0} (1 - |\Gamma_L|^2)$$

***Clearly**, this value is maximized when $\Gamma_L = 0$ (i.e., when $Z_L = Z_0$)!!!*

A: You are forgetting one very important fact! Although it is true that the load impedance Z_L affects the **reflected** wave power P^- , the value of Z_L —as we have shown in this handout—**likewise** helps determine the value of the **incident** wave (i.e., the value of P^+) as well.

- * Thus, the value of Z_L that minimizes P^- will **not** generally maximize P^+ !
- * **Likewise** the value of Z_L that maximizes P^+ will not generally minimize P^- .
- * Instead, the value of Z_L that maximizes the **absorbed** power P_{abs} is, by definition, the value that maximizes the **difference** $P^+ - P^-$.

We find that this impedance Z_L is the value that results in the **ideal** case of $Z_{in} = Z_g^*$.

Q: *Yes, but what about the case where $Z_g = Z_0$? For that case, we determined that the incident wave is independent of Z_L . Thus, it would seem that at least for that case, the **delivered** power would be maximized when the **reflected** power was minimized (i.e., $Z_L = Z_0$).*

A: True! But think about what the **input** impedance would be in that case— $Z_{in} = Z_0$. Oh by the way, that provides a **conjugate match** ($Z_{in} = Z_0 = Z_g^*$)!

Thus, in some ways, the case $Z_g = Z_0 = Z_L$ (i.e., both source and load impedances are numerically equal to Z_0) is **ideal**. A **conjugate match** occurs, the incident wave is **independent** of Z_L , there is **no** reflected wave, and all the math **simplifies** quite nicely:

$$V_0^+ = \frac{1}{2} V_g e^{-j\beta l} \qquad P_{abs} = P_{avl} = \frac{|V_g|^2}{8 Z_0}$$



Very Important Thing #2

Note the conjugate match criteria **says**:

Given source impedance Z_g , maximum power transfer occurs when the input impedance is set at value $Z_{in} = Z_g^$.*

It does **NOT** say:

Given input impedance Z_{in} , maximum power transfer occurs when the source impedance is set at value $Z_g = Z_{in}^$.*

This last statement is in fact **false!**

A **factual** statement is this:

Given input impedance Z_{in} , maximum power transfer occurs when the source impedance is set at value $Z_g = 0 - jX_{in}$ (i.e., $R_g = 0$).

Q: Huh??

A: Remember, the value of source impedance Z_g affects the available power P_{avl} of the source. To maximize P_{avl} , the real (resistive) component of the source impedance should be as small as possible (regardless of Z_{in} !), a fact that is **evident** when observing the expression for **available power**:

$$P_{avl} = \frac{1}{2} |V_g|^2 \frac{1}{4 \operatorname{Re}\{Z_g^*\}} = \frac{|V_g|^2}{8R_g}$$

Thus, **maximizing** the power delivered to a load (P_{abs}), from a source, has **two** components:

1. Maximize the **power available** (P_{avl}) from a source (e.g., minimize R_g).
2. **Extract** all of this available power by setting the input impedance Z_{in} to a value $Z_{in} = Z_g^*$ (thus $P_{abs} = P_{avl}$).